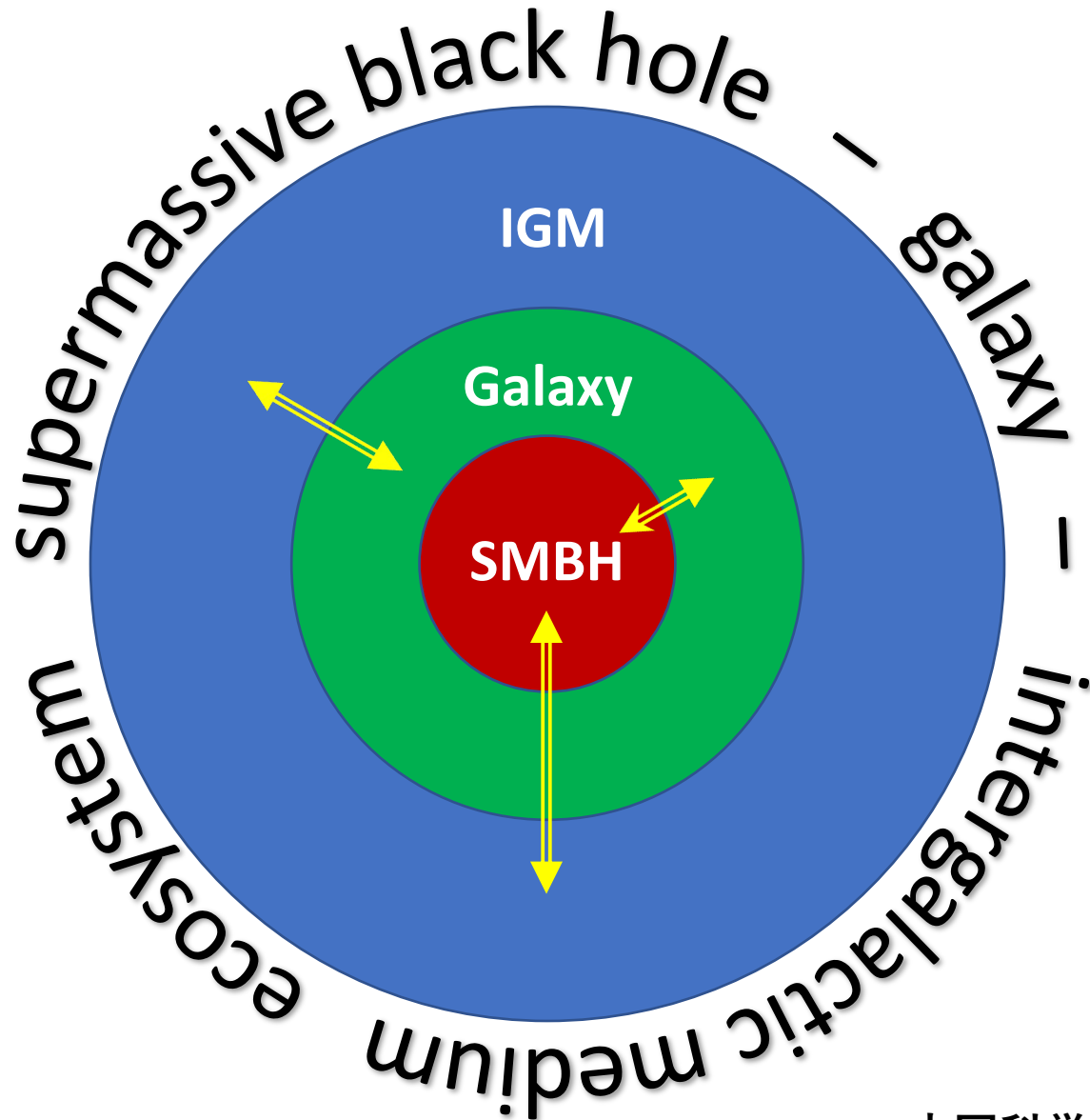


Physical Mechanisms Regulating Gas Supply to Galaxies



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&

Institute of Astronomy

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@中国科学院国家天文台, May 17, 2023



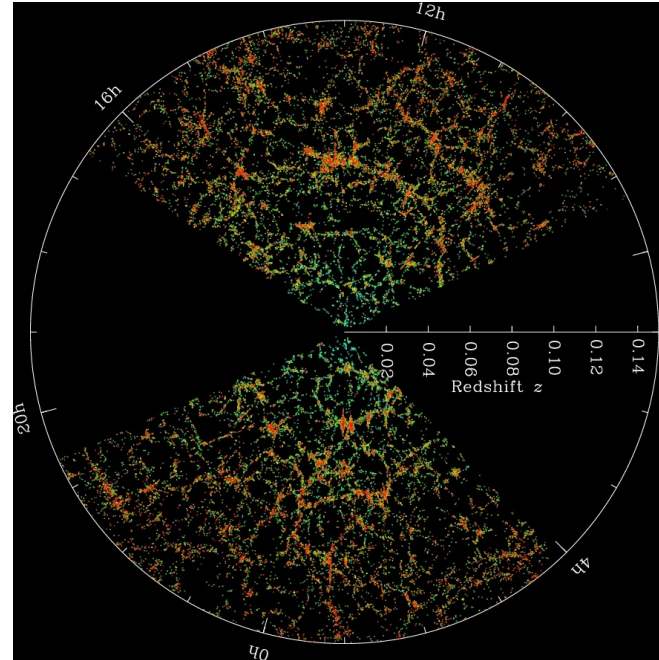
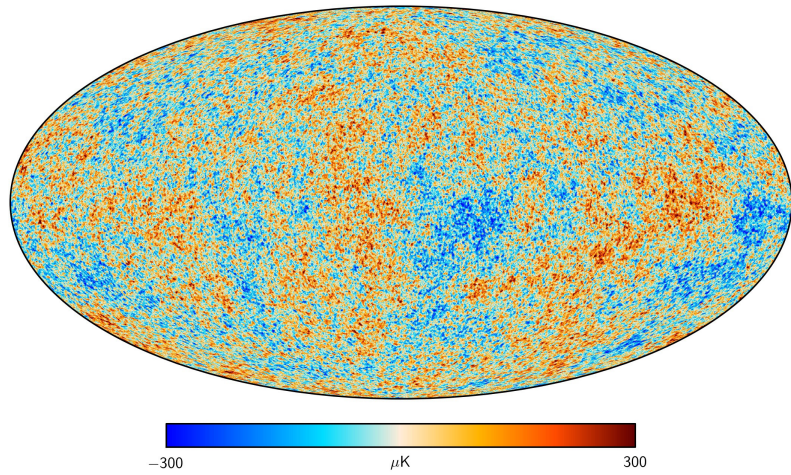
Outline

- **The core astrophysical problem**
- **A multi-scale computational problem**
- **A multi-physics problem: cosmology, dark energy, dark matter, gravity, gas dynamics, microphysics, chemistry, stellar feedback, supermassive black hole feedback**
- **Conclusions**



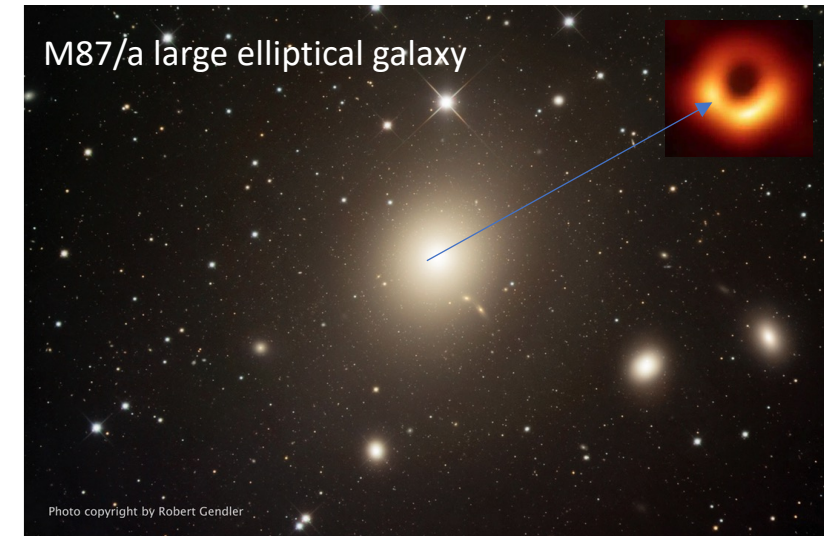
The core astrophysical problem : **how are galaxies formed?**

- **goal: learn physics of galaxy formation**



Planck cosmic microwave background surveyor

SDSS local universe



Outline

- **The core astrophysical problem**
- **A multi-scale computational problem**
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The core astrophysical problem : **A hard computational problem**

Multi scales

- **star forming clouds: 0.1pc**
- **galaxies: 100kpc**
- **intergalactic medium: 100 Mpc**
- **universe: 10^4 Mpc**

Scales are dynamically coupled

- **gravity is of long range**
- **structure seed fluctuation energy peaks at 180Mpc**
- **radiation and cosmic rays act over very large scales**
- **atomic scale processes influence how galaxies form**
- **galactic feedback reaches beyond galactic scales**

11 dex computational symphony



Five pianos worth of dynamic range that are coupled,

Outline

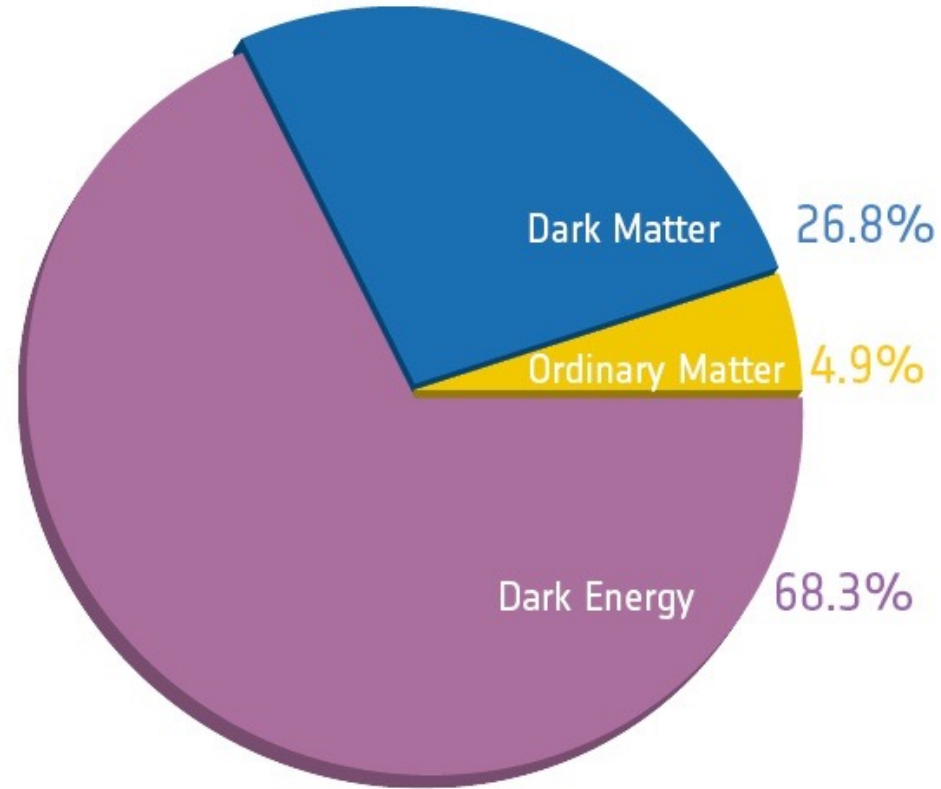
- **The core astrophysical problem**
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- **Conclusions**



The core astrophysical problem : **A multi-physics problem**



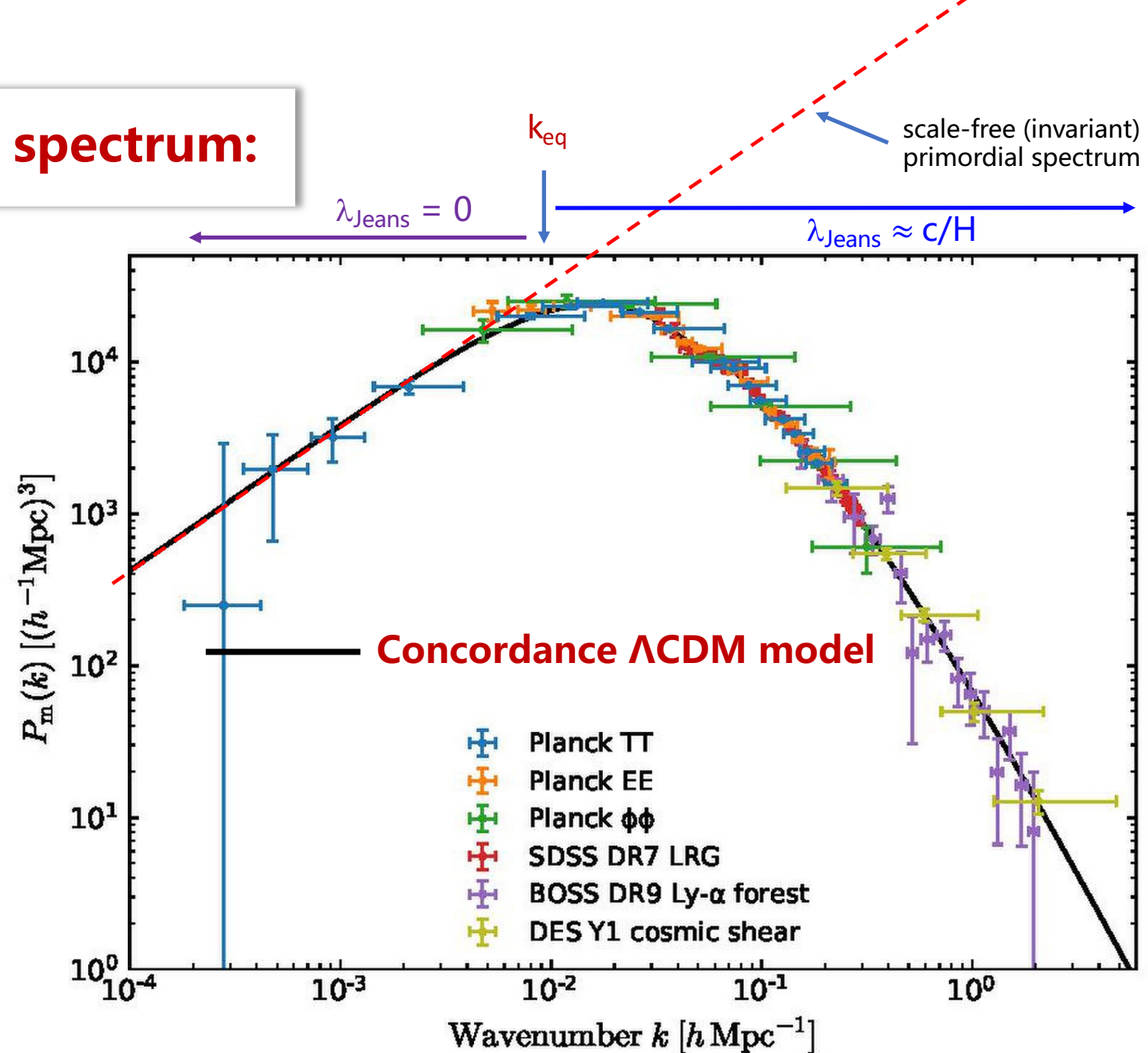
A flat Λ -dominated cold dark matter universe:



Take-away message: consistent with inflation
Puzzle: partitioning of energy density is mystery

Matter power spectrum:

- Baryon-photon is **dynamically coupled** to cold dark matter (CDM) by gravity prior to recombination, so CMB power spectrum yields CDM power spectrum, $P_m(k)$
- The scale-free primordial cold dark matter power spectrum, generated by quantum fluctuations prior to inflation, is later modified by the transfer function, i.e. **Jeans suppression**
- In radiation dominated (redshift > 3400) era, total matter is relativistic and Jeans length is **equal to the horizon scale**, while in matter dominated era, Jeans length **drops to zero**
- Fourier modes that enter the horizon after radiation-matter equality (**eq**) redshift are unchanged, while earlier, shorter modes are suppressed



Hubble constant crisis (?) : 67.4 vs 73.3 km/s/Mpc

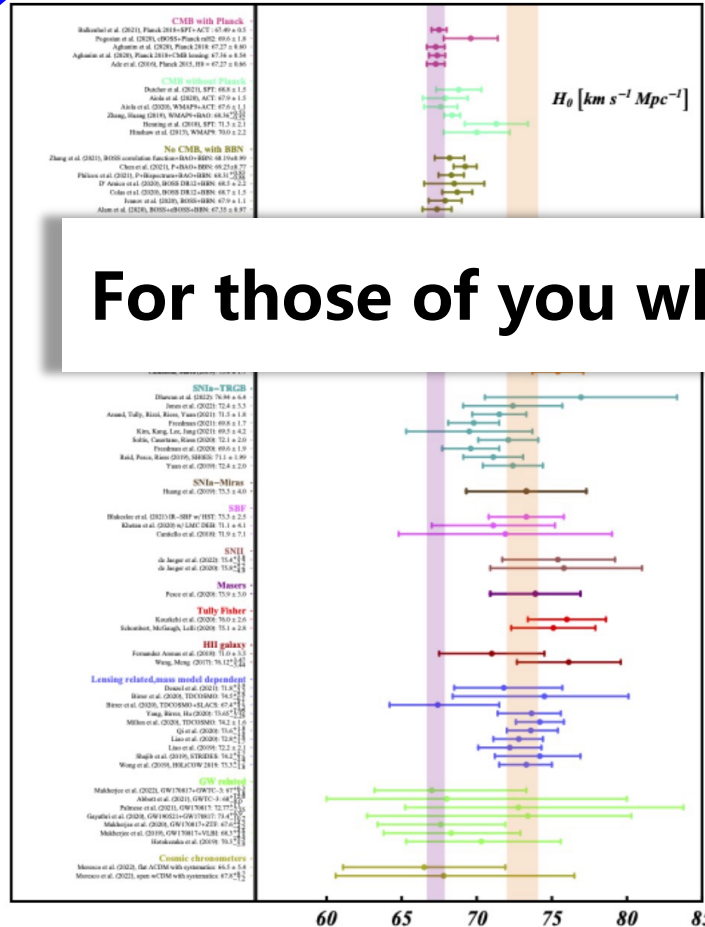


Figure 2

68% confidence-level constraints on H_0 from different cosmological probes. From Ref. (12) (based on Refs. [49, 50]).

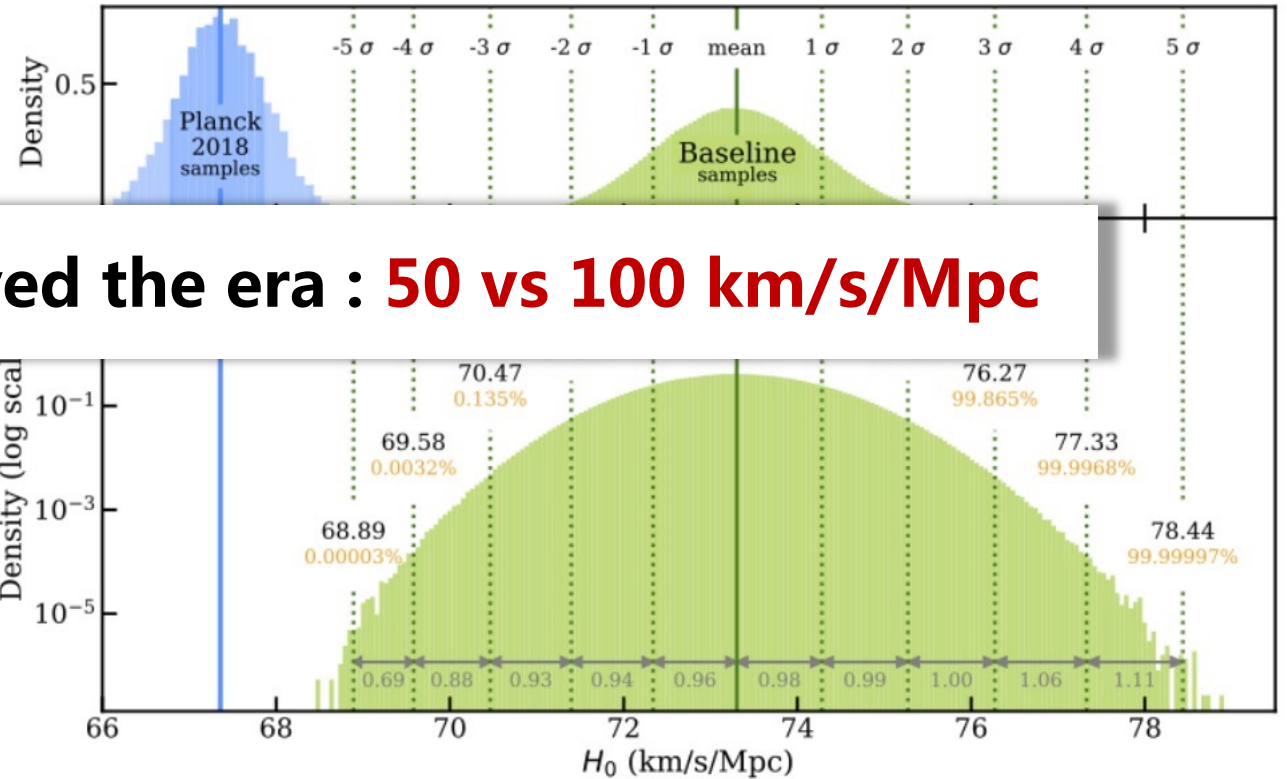


Figure 1

Extended MCMC sampling of the posterior for H_0 to measure out to the 5σ confidence level. The upper panel shows the probability density for the baseline from SH0ES and from the Planck Collaboration et al. (2020) chains. The bottom panel shows the log of the probability density to improve the ability to see the tails.

dark
energy

In 1915, Einstein gave the world his
General Theory of Relativity

Einstein tensor:
curvature of spacetime

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu}$$

Einstein added the cosmological
constant, Λ , **the anti-gravity**.
It was the biggest blunder of his life,
said he (according to Gamow).



Albert Einstein, Hideki Yukawa, and John Archibald Wheeler walk through Marquand Park in Princeton, New Jersey, in 1954.

or was it?

Understanding origin of Λ may
shed light on new physics to the
standard particle physics model

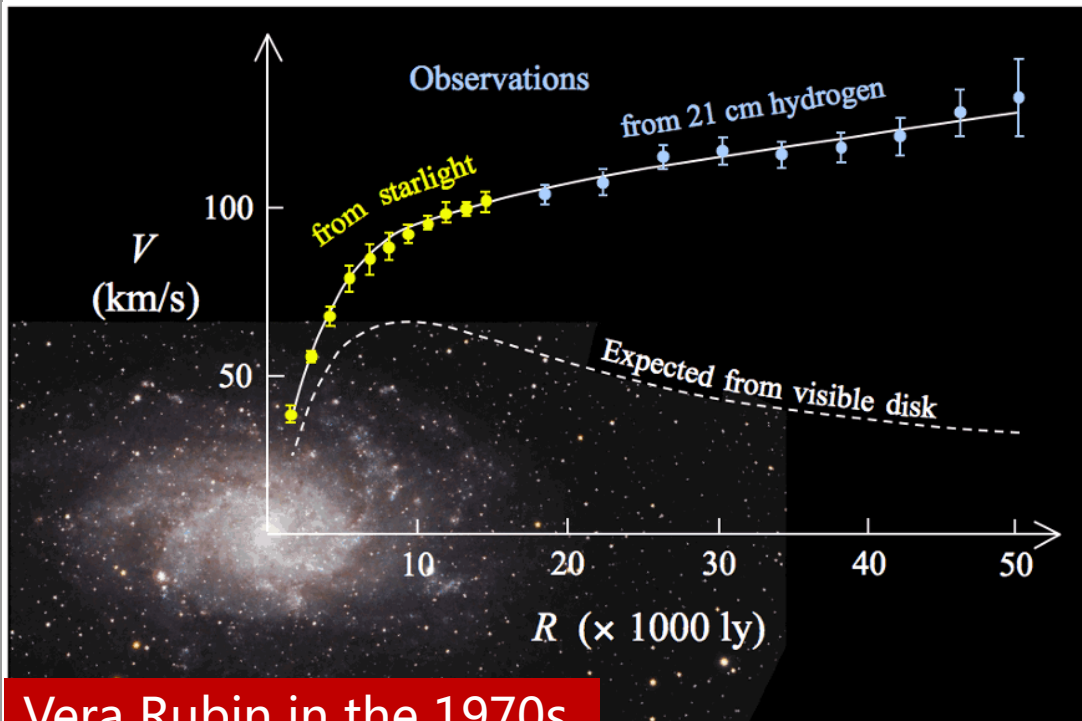


dark
matter

Two-fold evidence (implied, earlier)

Present-day galaxies, clusters:

➤ Rotation curves of stars and gas



Vera Rubin in the 1970s

Galaxy formation:

➤ DM to seed galaxy formation

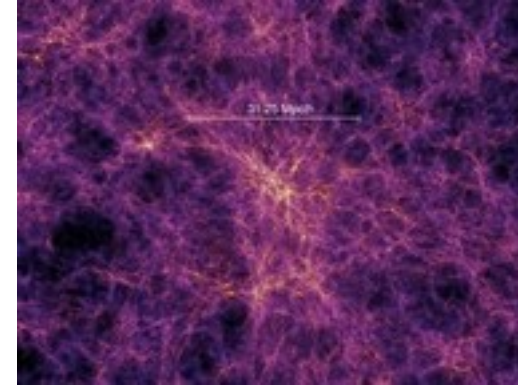
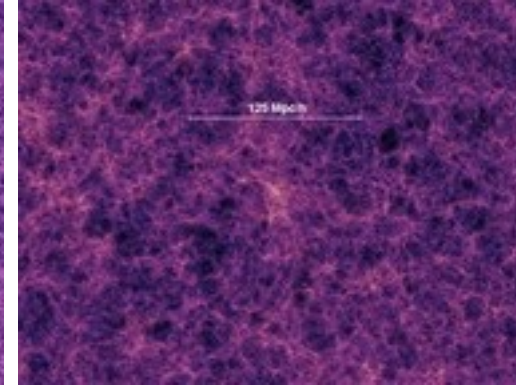
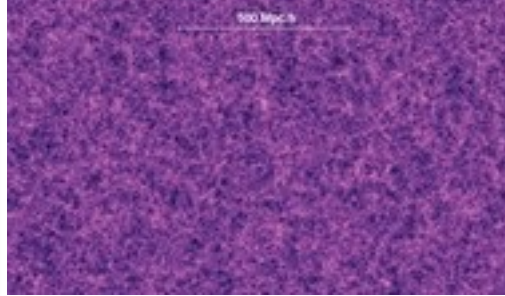
- $z > 1100$ the photon-baryon fluid is only **nearly perfect**
- **On small scales**, the coupling is **not perfect** and baryons and photons can past one another (Silk 1968) at $z > 1100$, i.e., **adiabatic** perturbations with wavelengths shorter than photon mean free path are damped
- It implies that dark matter is needed to **seed** galaxy formation after recombination
- Otherwise, galaxies **less massive than $10^{11} M_{\odot}$** **will not exist** in most places, but we see them everywhere

gravity

Dynamics of dark matter after 40 years of effort

- We understand the gravitational **growth** of dark matter from recombination to the present
- We understand the **growth**, **abundance** and **clustering** of dark matter halos
- We understand that the density profile of virialized dark matter halos decreases monotonically towards larger radii from center, although exact slope at the very center is not universally agreed upon and variations exist

Redshift $z=5.7$
 $t = 1$ Gyr



Redshift $z=1.4$
 $t = 4.7$ Gyr



Redshift $z=0$
 $t = 13.6$ Gyr



Millennium Simulations (Springel et al 2005)

Basic structure formation picture:

- Structure formation proceeds **in a hierarchical fashion**: small scales become nonlinear and collapse first
- Earliest structures of mass **$10^5 M_\odot$** form at redshift 20-30
- Mergers and accretion of smaller structures → ever larger structures with time
- Cycle has gone on, will cease soon due to increasingly dominant positive Λ energy
- The largest gravitationally bound systems are **clusters of galaxies** of size 10 million light years across (probing the 2σ tail of the Gaussian distribution, i.e., σ_8, Ω_M)
- Looking ahead, the Milky Way and the Andromeda may collide in about **6 billion years**, and Λ cannot stop it, if you were counting on Λ

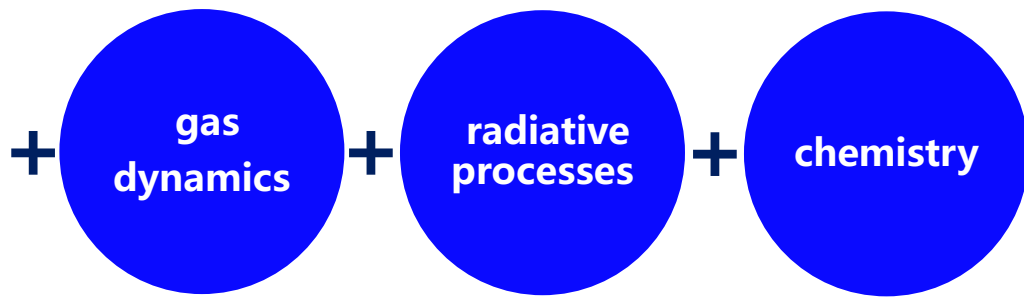
What about galaxies ?

**It is mainly about cold gas supply,
we need to understand both**

internal and external

**physical processes that ultimately
affect the cold gas supply to
galaxies**





gas processes linked to gravity

**Two robust results from cosmological hydrodynamic simulations mark
outer boundary conditions of galaxies mainly influenced by gravity**

Warm-hot intergalactic medium

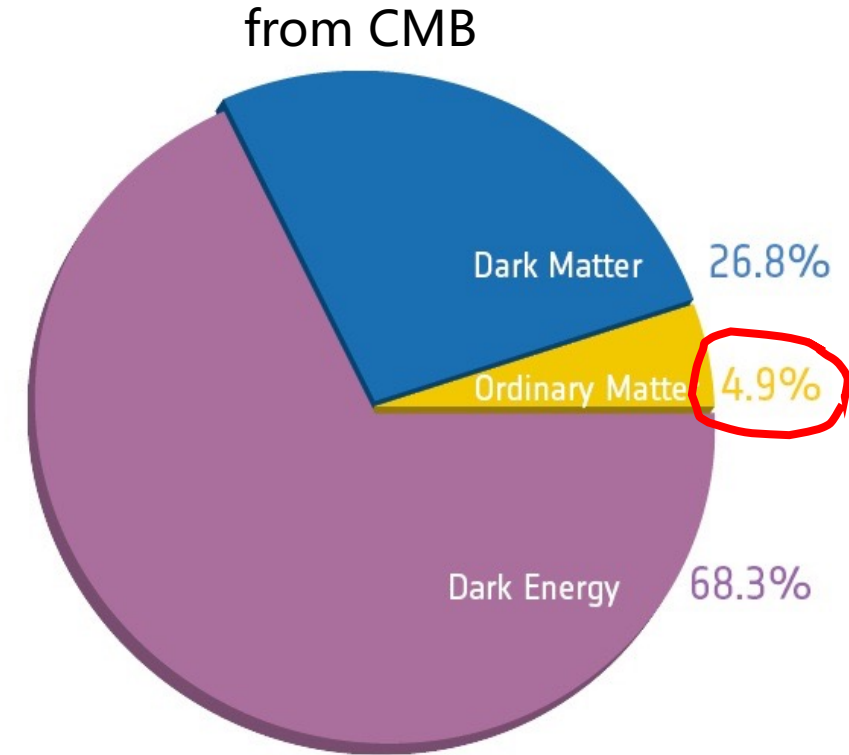
hot-cold accretion dichotomy



TABLE 3
THE BARYON BUDGET

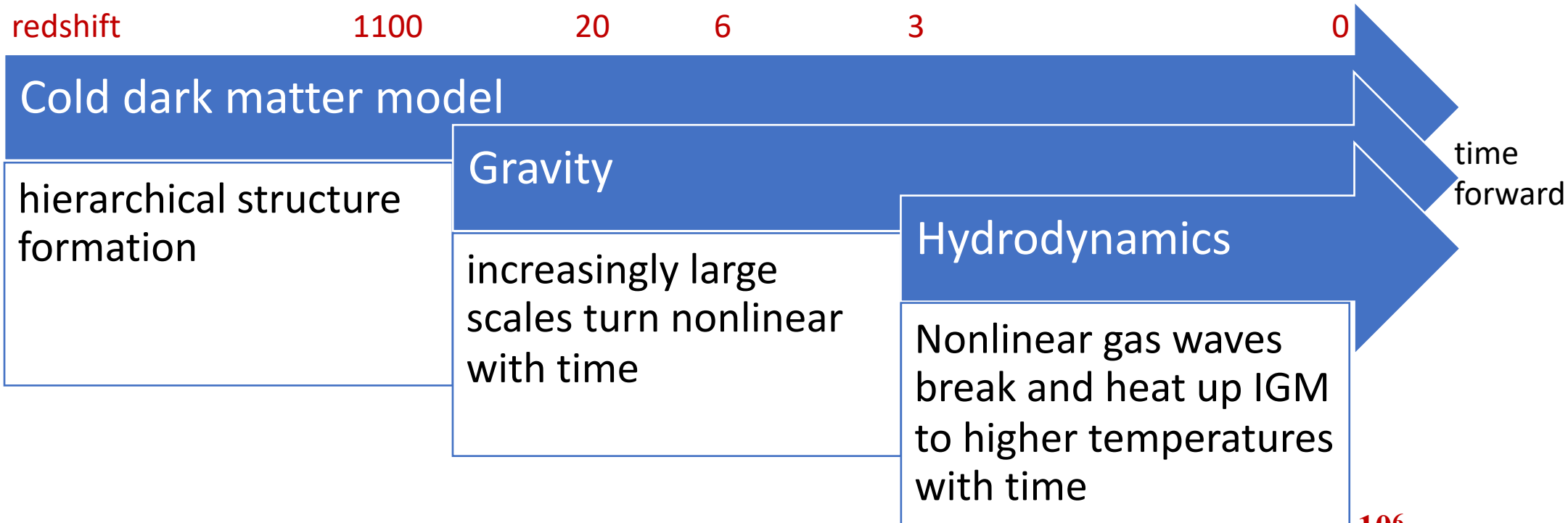
Component	Central	Maximum	Minimum	Grade ^a
Observed at $z \approx 0$				
1. Stars in spheroids	$0.0026 h_{70}^{-1}$	$0.0043 h_{70}^{-1}$	$0.0014 h_{70}^{-1}$	A
2. Stars in disks	$0.00086 h_{70}^{-1}$	$0.00129 h_{70}^{-1}$	$0.00051 h_{70}^{-1}$	A—
3. Stars in irregulars	$0.000069 h_{70}^{-1}$	$0.000116 h_{70}^{-1}$	$0.000033 h_{70}^{-1}$	B
4. Neutral atomic gas	$0.00033 h_{70}^{-1}$	$0.00041 h_{70}^{-1}$	$0.00025 h_{70}^{-1}$	A
5. Molecular gas	$0.00030 h_{70}^{-1}$	$0.00037 h_{70}^{-1}$	$0.00023 h_{70}^{-1}$	A—
6. Plasma in clusters	$0.0026 h_{70}^{-1.5}$	$0.0044 h_{70}^{-1.5}$	$0.0014 h_{70}^{-1.5}$	A
7a. Warm plasma in groups	$0.0056 h_{70}^{-1.5}$	$0.0115 h_{70}^{-1.5}$	$0.0029 h_{70}^{-1.5}$	B
7b. Cool plasma	$0.002 h_{70}^{-1}$	$0.003 h_{70}^{-1}$	$0.0007 h_{70}^{-1}$	C
7'. Plasma in groups	$0.014 h_{70}^{-1}$	$0.030 h_{70}^{-1}$	$0.0072 h_{70}^{-1}$	B
8. Sum (at $h = 70$ and $z \approx 0$)	0.021	0.041	0.007	...
Gas components at $z \approx 3$				
9. Damped absorbers	$0.0015 h_{70}^{-1}$	$0.0027 h_{70}^{-1}$	$0.0007 h_{70}^{-1}$	A—
10. Ly α forest clouds	$0.04 h_{70}^{-1.5}$	$0.05 h_{70}^{-1.5}$	$0.01 h_{70}^{-1.5}$	B
11. Intercloud gas (He II)	$0.01 h_{70}^{-1.5}$	$0.0001 h_{70}^{-1}$	B
Abundances of:				
12. Deuterium	$0.04 h_{70}^{-2}$	$0.054 h_{70}^{-2}$	$0.013 h_{70}^{-2}$	A
13. Helium	$0.010 h_{70}^{-2}$	$0.027 h_{70}^{-2}$...	A
14. Nucleosynthesis	$0.020 h_{70}^{-2}$	$0.027 h_{70}^{-2}$	$0.013 h_{70}^{-2}$...

^a Confidence of evaluation, from A (robust) to C (highly uncertain).



Where are (the other half of) the cosmic baryons ?

Physics



temperature of intergalactic medium [Kelvin]

10^{12}

$10^{3.5}$

10^0

10^4

reionization

photo heating

Shock heating: $T_{\text{IGM}} \propto (H(z) L_{\text{nonlinear}})^2$

10^6

Predictions

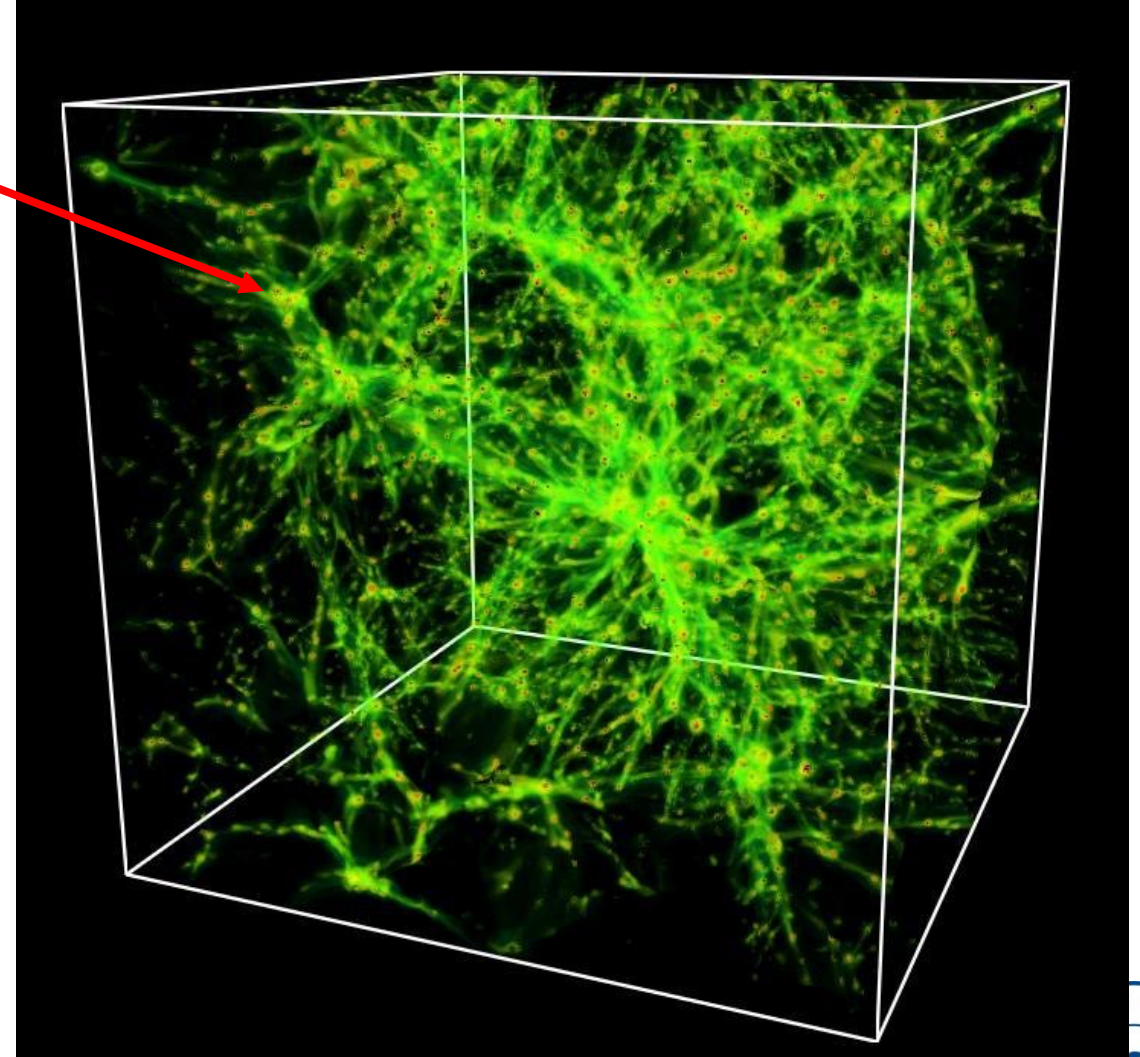
- 50% of cosmic baryons in the universe @ $z=0$ exist in a filamentary network with temperature of 10^5 - 10^7 K: **the warm-hot intergalactic medium (WHIM)** (Cen & Ostriker 1999, ApJ, 514, 1) :

Confirmation

- Warm component with $T = 10^5$ - 10^6 K was unambiguously confirmed by the **Hubble Space Telescope** COS far UV absorption line observations

on galaxy formation

- **WHIM at $z=0$ -3 poses as an important boundary condition for galaxies accreting from it**



Warm-hot intergalactic medium

: further confirmation imperative

Further confirmation
in X-ray is needed to
complete the checking

X棱镜

XRISM X-Ray Imaging and Spectroscopy Mission

Launch time : 2023

XRISM : jointly developed by JAXA/NASA (日本宇宙航空研究开发机构和NASA)

Resolving the Nature of the Energetic Cosmos

XRISM major science subjects include:
The missing baryons in the intergalactic medium

3.1 Missing Baryons and the WHIM

Current observations of hot gas in galaxy clusters, the Local Group, and the warm intergalactic medium (IGM) indicate that 20-40% of the baryons seen in the high- z Universe are missing from our observations at $z=0$. The missing baryons are suspected to occupy the filaments of gas

清华大学正开展的项目

HUBS: 热宇宙重子测量员

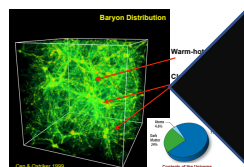
HUBS: Hot Universe Baryon Surveyor

Wei Cui, for the HUBS Collaboration
Department of Astronomy, Tsinghua University



INTRODUCTION

Based on the measured abundances of primordial isotopes the theory of Big Bang Nucleosynthesis (BBN) tells us how many baryons were produced nearly at the birth of the universe. As the universe evolved, it cooled and, in a matter of minutes after the Big Bang, it could no longer sustain fusion, so the production of heavier isotopes ends. As it cooled further, neutral atoms began to form about 370,000 years after the Big Bang, locking up the electrons that had tied photons and matter together via scattering and allowing the photons to fill the entire universe, which produced what is now known as the cosmic microwave background (CMB). Precise measurements on the anisotropy of the CMB have allowed us to derive the values of key cosmological parameters, including the energy density of baryonic matter. The result implies that all of the BBN baryons were still there at the redshift of about 1100. Going down in redshifts, first stars and galaxies formed and evolved, when the baryons seem to go "missing" in optical surveys. In the present-day universe, only about half of the baryons are seen optically. This is the long-standing "missing baryon problem".



Theoretically, cosmological hydrodynamical simulations show a significant fraction of cosmic baryons were heated to temperatures of 10^4 K by shocks that had been produced during the formation of large-scale structures. Radiation from such low-density, hot baryons lies mainly between far UV and soft X-ray wavelengths, making it difficult to detect observationally. Such a solution to the "missing baryon problem" has found significant support from indirect observations, through the absorption of radiation from distant quasars by the missing baryons or their distortion of the CMB spectrum.

The primary scientific objectives of HUBS include: (1) to directly detect X-ray emission from the hot baryons in the IGM or CGM, and characterize their physical and chemical properties; and (2) to study, based on the observations, the accretion and feedback processes that are thought to be highly relevant to the heating and chemical enrichment of the baryons in the CGM (and perhaps also IGM). The results are expected to help advance our understanding of galaxy formation and evolution significantly. Secondary objectives are many, including hot interstellar medium, diffuse X-ray background, supernova remnants, as well as charge exchange processes in the solar system.

PRELIMINARY DESIGN

The most effective approach is to carry out high-throughput, high-resolution imaging spectroscopy in the soft X-ray band (< 2 keV), where the hot baryons are expected to manifest themselves in emission and absorption lines. The best spectral lines to focus on are associated with helium-like and hydrogen-like ions (especially oxygen lines at ~ 0.6 keV).

The HUBS imaging spectrometer will be based on the transition-edge sensor (TES) technology that is optimized for the softest X-ray band (0.1-2 keV), where the oxygen lines are located. The design aims at maximizing the product of grasp ($A_{\text{eff}} \Omega$) and spectral resolution ($E/\Delta E$), for detecting extended X-ray emission,² and achieving moderate angular resolution at the same time.

ENABLING TECHNOLOGIES

The enabling technologies of HUBS include: TES array with large X-ray absorbers, SQUID-based multiplexing readout, cryocooler and ADR, and X-ray optics.

A 2-stage cryocooler design is baselined for providing a pre-cooled stage for the ADR. Two prototypes have been developed³⁻⁵, with one referred to as Vuilleumier hybrid pulse tube cryocooler (VM-PTC), and the other High-frequency pulse tube cryocooler (HPTC). VM-PTC consists of a VM stage and a PT stage, and is able to reach 2.17 K, providing about 10 mW of cooling power. The HPTC consists of two PT stages, and is able to reach 3.5 K, providing about 6 mW of cooling power.

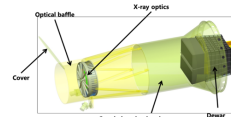
As for X-ray optics, the slumped glass technique is used to form light-weight, thin glass mirror shells, and the shells are then nested to realize Wolter-I optics with a large

清华大学主持HUBS项目，
拟用X发射线来绘制出这宇宙网络，来检测预言的**热星系际介质**，由中国国家航天局资助，现处关键技术开发阶段

Launch time : 2030

The grasp of HUBS is, therefore, over an order of magnitude larger than Athena X-IFU, and is, by design, complementary in primary scientific objectives.

HUBS adopts an integrated design of the scientific payload and the satellite platform.



HUBS intends to sustain the scientific community's quest for a dedicated X-ray mission to detect the missing baryons over the past nearly two decades. At present, science definition and technology R&D are mainly being carried out by collaborating universities and research institutes in China, but international collaboration is strongly encouraged, by CNSA, at all levels, including observing strategies and target selection through the Science Working Groups, technology R&D, and auxiliary payloads (which enhance and/or broaden the scientific capabilities and objectives of HUBS).

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1. R. Cen, & J.P. Ostriker, *Astrophys. J.*, 515, L109 (1999)
2. T. Fang, et al., *Astrophys. J.*, 625, 612 (2005)
3. J. Wang, C. Fan, T. Zhang, et al., *Sci. Bull.*, 64, 219 (2019)
4. C. Fan, J. Wang, K. Liu, et al., *Cryogenics*, 98, 71 (2019)
5. L. Chen, X. Wu, J. Wang, et al., *Cryogenics*, 94, 103 (2018)

Athena: ESA's Big X-ray Telescope

By Elizabeth Howell



雅典娜: 欧洲航天局的大X射线望远镜

Launch time : 2035



Artist's conception of the proposed Advanced Telescope for High-Energy Astrophysics (Athena+). (Image credit: Athena+ Team)

The spacecraft is an X-ray observatory expected to answer questions about the hot and energetic universe. The two main questions include "How does ordinary matter assemble into the large-scale structures that we see today" (such as galaxy groups and clusters) and "How do black holes grow and influence the environment around them." A black hole's environment is best seen in X-rays, and hot gases, which are visible in X-rays, can give clues about galactic structure

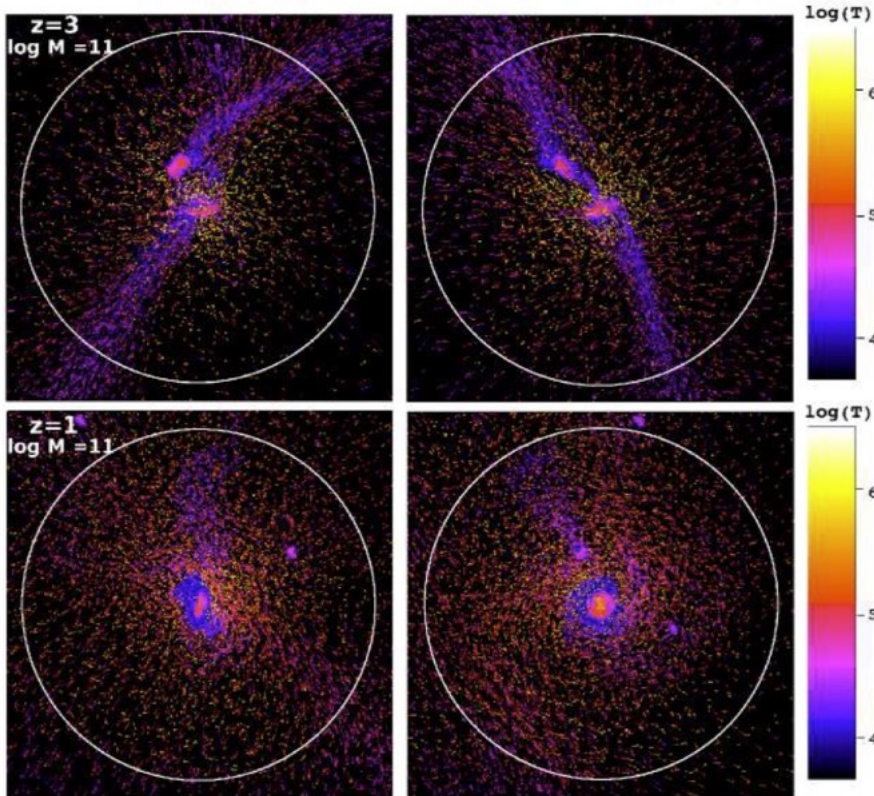
One of the main science subjects:
The missing baryons in the intergalactic medium

- missing baryons in the intergalactic medium, or plasma between galaxies;
- how black holes and supermassive black holes grow over time;

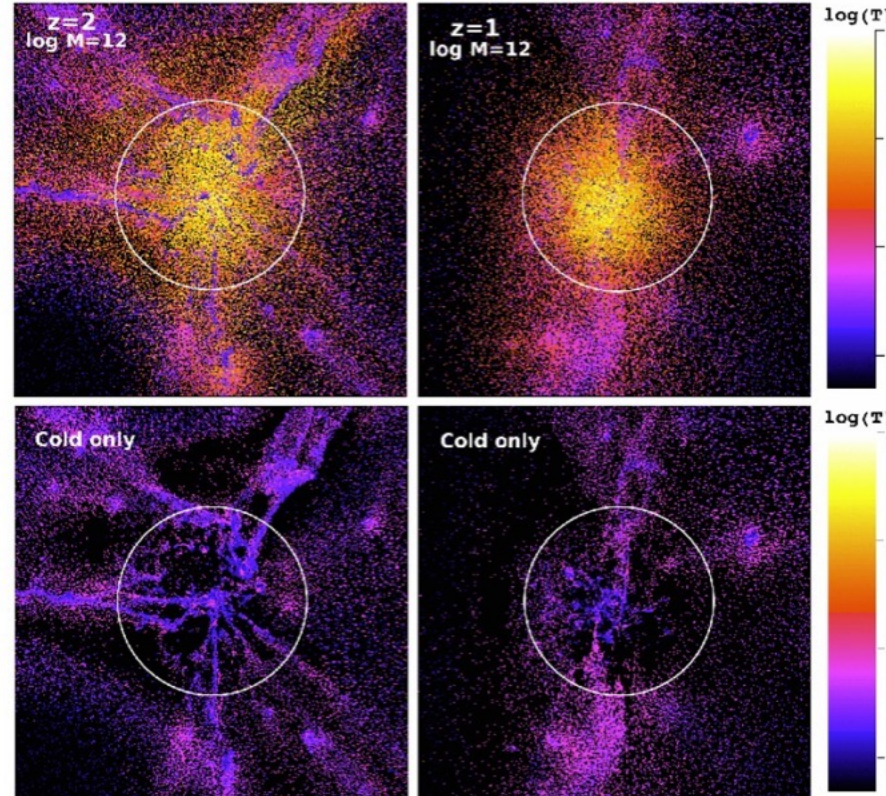
For more information, please visit the HUBS web site: <http://hubs.phys.tsinghua.edu.cn/en/>



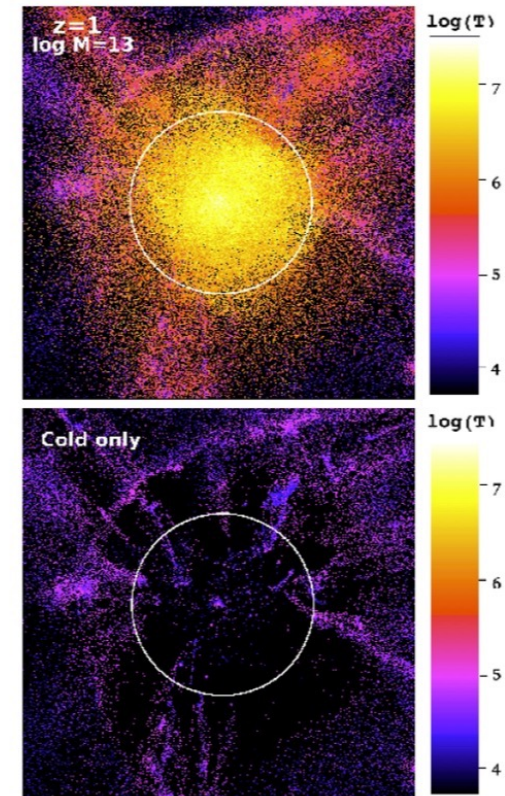
$M_h = 10^{11} M_\odot$



$M_h = 10^{12} M_\odot$

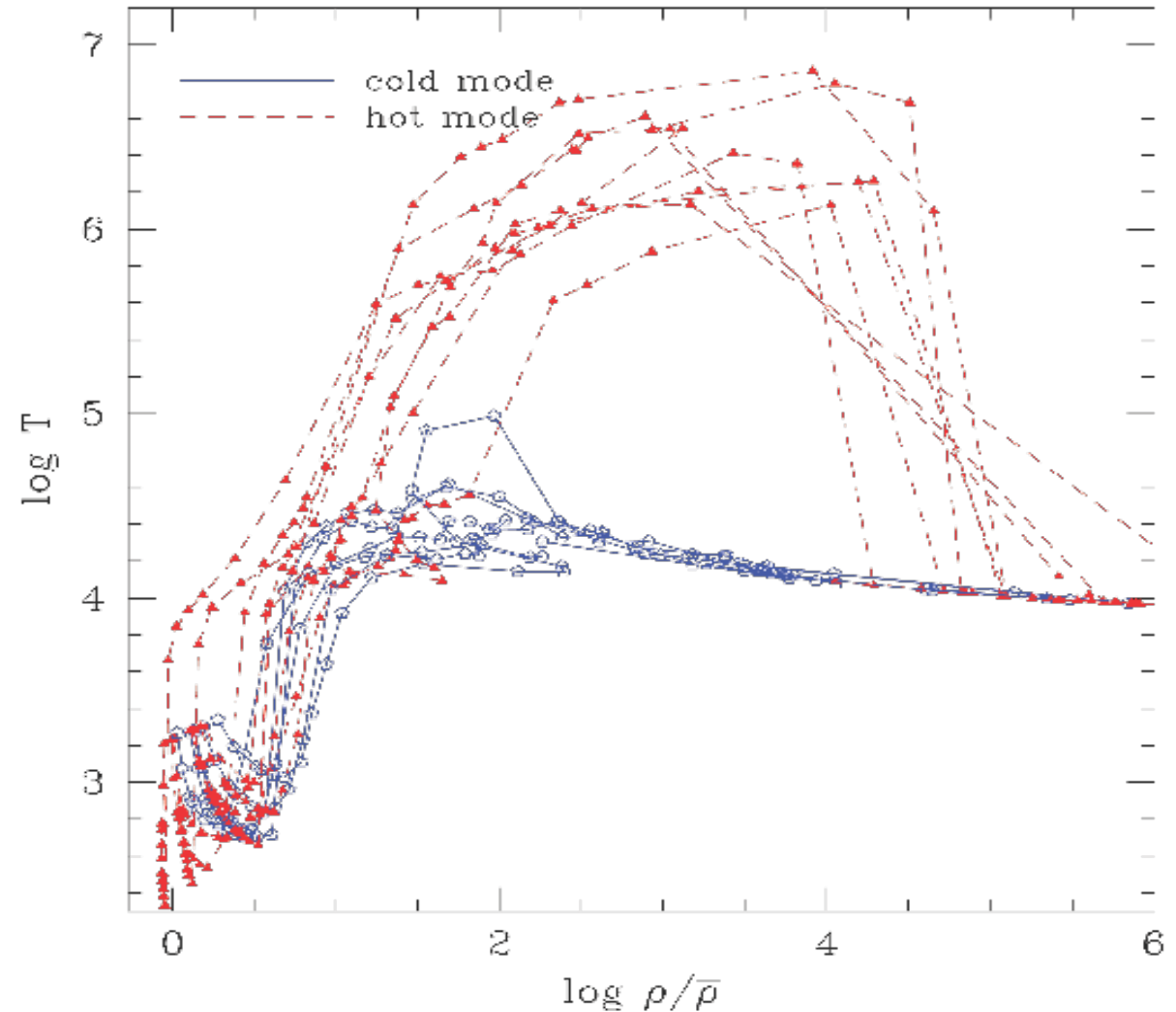


$M_h = 10^{13} M_\odot$



Significance

- Classic picture: all gas entering halo is shocked heated to virial temperature
- New insight: not all accreted gas is shock heated to virial temperature entering halo
- There are **two distinct modes: cold and hot accretion modes**
- Halo mass dependent accretion mode has major implications for galaxy formation

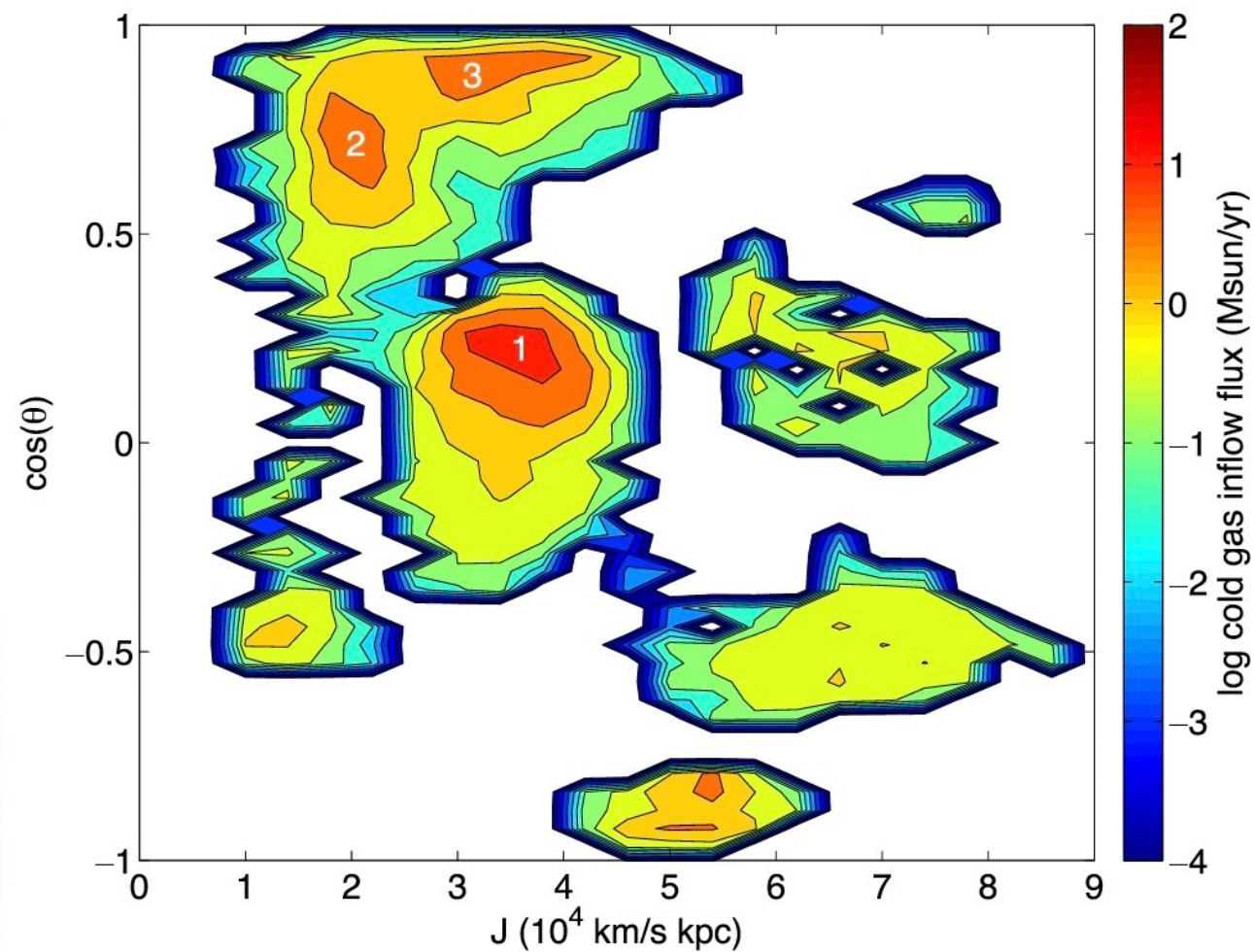
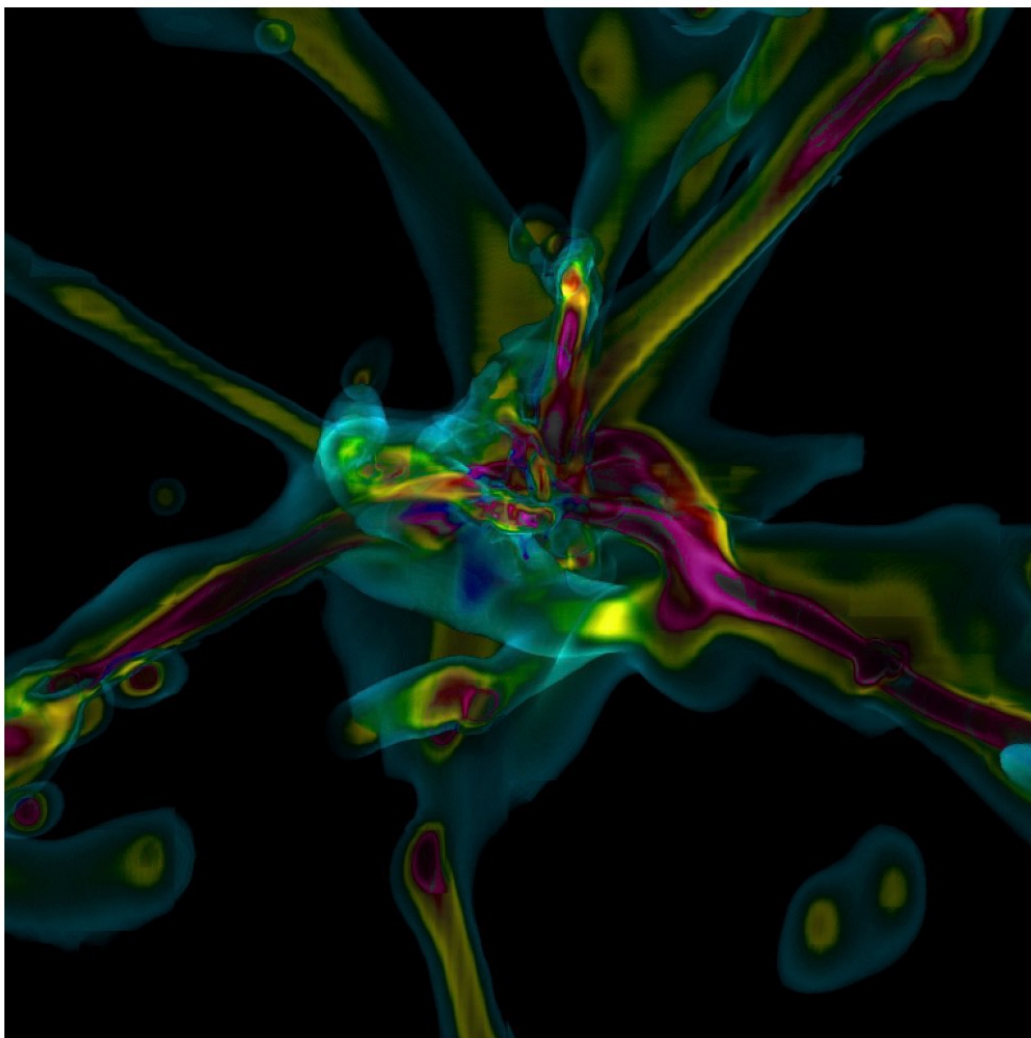


Keres et al 2005, MNRAS, 363, 2

hot-cold accretion dichotomy

: quantifying cold streams

cold streams are readily identifiable in cosmological hydrodynamic simulations



Stream interactions : **SMBH mass to bulge mass ratio**

mass_{gas}(J) of episodic inflow could lead to $M_{\text{BH}}-M_{\text{bulge}}$ relation

- Simulations (Hopkins & Quataert 2010,2011) show that from 0.01-1000pc inflow gas ends with a Mestel disk of surface density $\Sigma(r) \propto r^{-1}$ (r is radius)

- $$r_0 = 0.42 (\alpha/0.1)^{2/5} (l_E/0.1)^{-2/5} (M_{\text{BH}}/10^8 M_\odot)^{3/25} \times (Ma/0.1)^{14/25} (\kappa/\kappa_e)^{4/25} \text{ pc} \quad (7)$$

Toomre $Q < 1$ for $r > r_0$

$Q > 1$ for $r < r_0$

(Goodman 2003, Equation (42)), where α is radiative efficiency, l_E is luminosity in Eddington units, Ma is Mach number of the viscous disk at r_0 , and κ and κ_e opacity, and electron-scattering opacity, respectively. Hence, the feeding rate to the accretion disk

- we combine the above two ingredients with 100pc resolution cosmological simulations to show that the observed $M_{\text{BH}}-M_{\text{bulge}}$ relation can be produced (Cen 2015)



$M_{\text{BH}}\text{-}M_{\text{bulge}}$ relation : reproduced from inflow gas J distribution

THE ASTROPHYSICAL JOURNAL LETTERS, 805:L9 (6pp), 2015 May 20

CEN

Blue: \log_{10} (accretion rate \dot{M}_{dot})
Red: \log_{10} ($500\dot{M}_{\text{dot}}/\text{SFR}$)

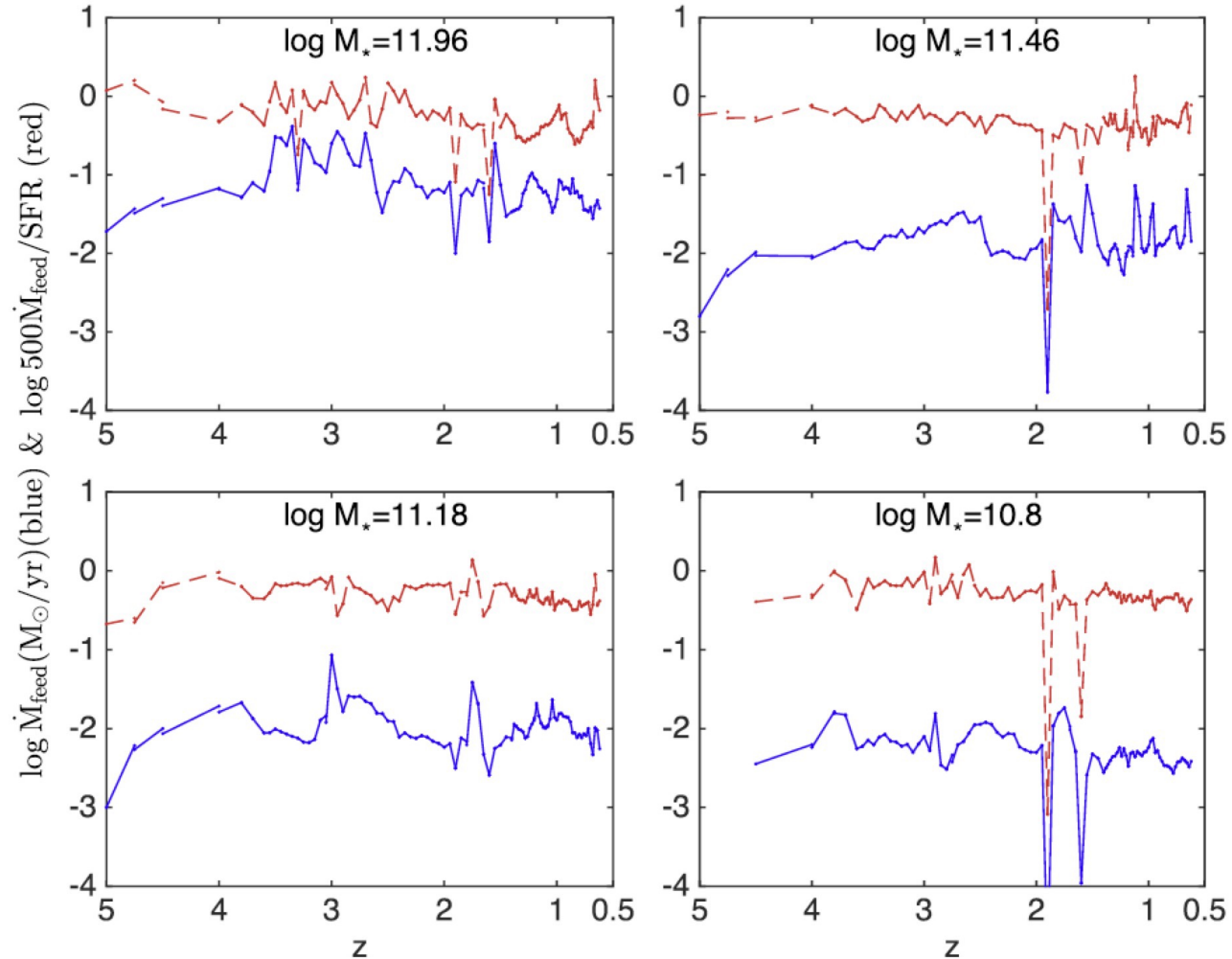


Figure 1. Shows histories of the feeding rate \dot{M}_{feed} (blue) and $R \equiv 500 \dot{M}_{\text{feed}} / \text{SFR}$ (red) for four random galaxies. The logarithm of the stellar mass for each galaxy at $z = 0.62$ is indicated at the top of each panel.

$M_{\text{BH}}-M_{\text{bulge}}$ due to inflow **regulation** : $M_{\text{BH}}/M_{\text{bulge}}$ increases with redshift

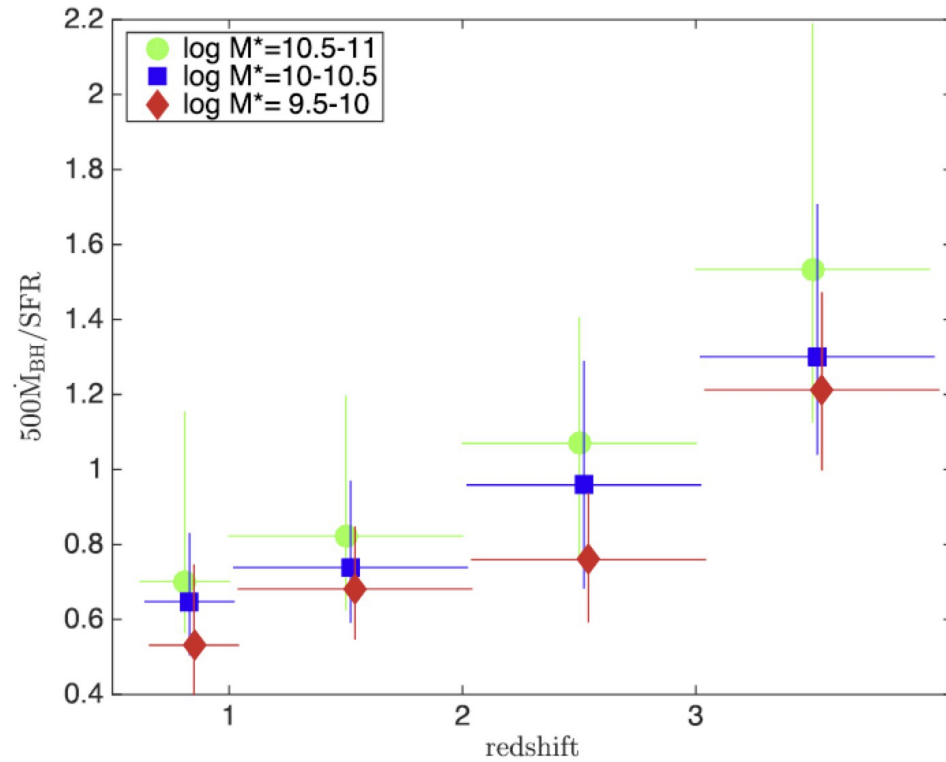


FIG. 12.— Median of R as a function of redshift, separately for three stellar mass ranges $10^{9.5-10} M_{\odot}$ (red), $10^{10-10.5} M_{\odot}$ (blue), and $10^{10.5-11} M_{\odot}$ (green), measured at the redshift in question. The vertical error bars indicate the interquartile range, whereas the horizontal error bars represent the uncertainty in redshift. The red and blue points are horizontally slightly right-shifted for clarity of display. There are (659, 2214) galaxies with stellar mass in the ranges $10^{9.5-10} M_{\odot}$ and $10^{10-10.5} M_{\odot}$, respectively.

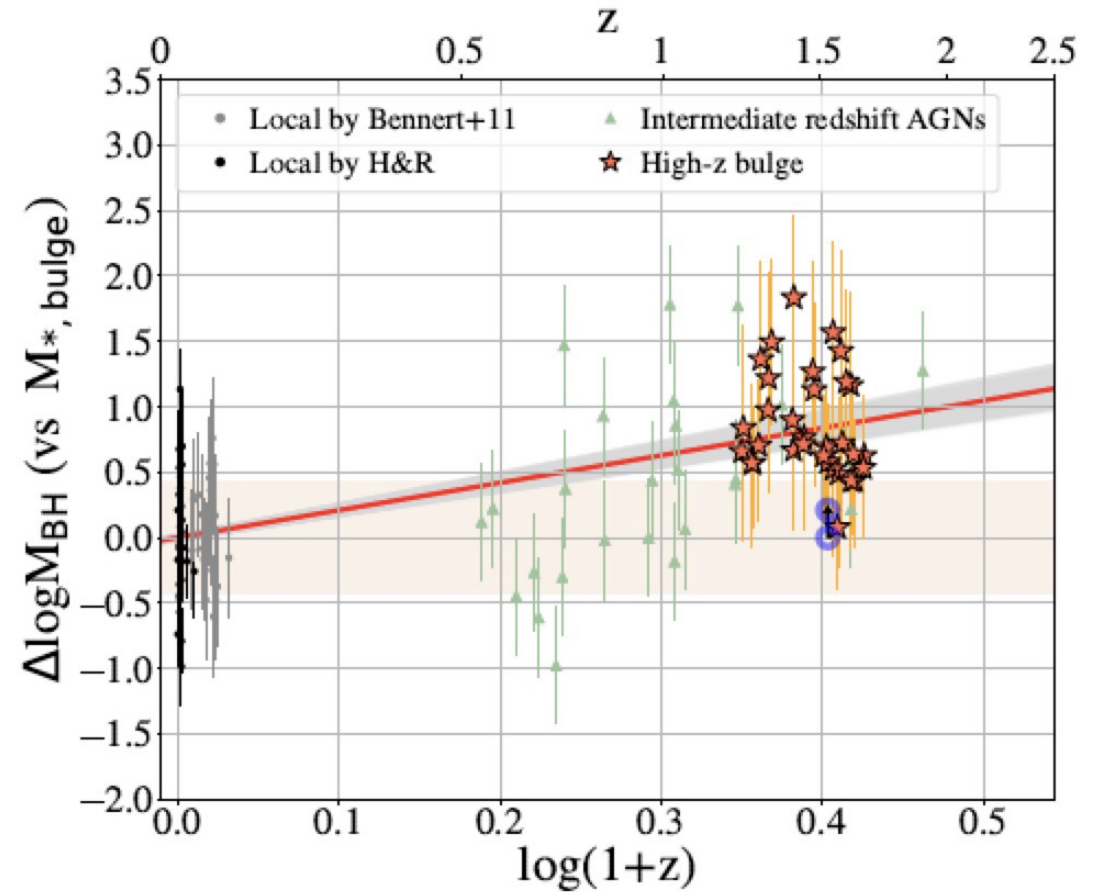


FIG. 13.— Same as Figure 8, but for $M_{\text{BH}}-M_{*, \text{bulge}}$ relations. The intermediate redshift AGNs are from Silverman (2013), listed in Table 7. Cisternas et al. (2011) do not provide information on the bulge mass.

$M_{\text{BH}}-M_{\text{bulge}}$ due to inflow regulation: enhanced central SF in QSOs

THE ASTROPHYSICAL JOURNAL, 944:30 (21pp), 2023 February 10

Molina et al.

Latest detailed observation shows significant central star formation enhancements in PG quasars, not the way around

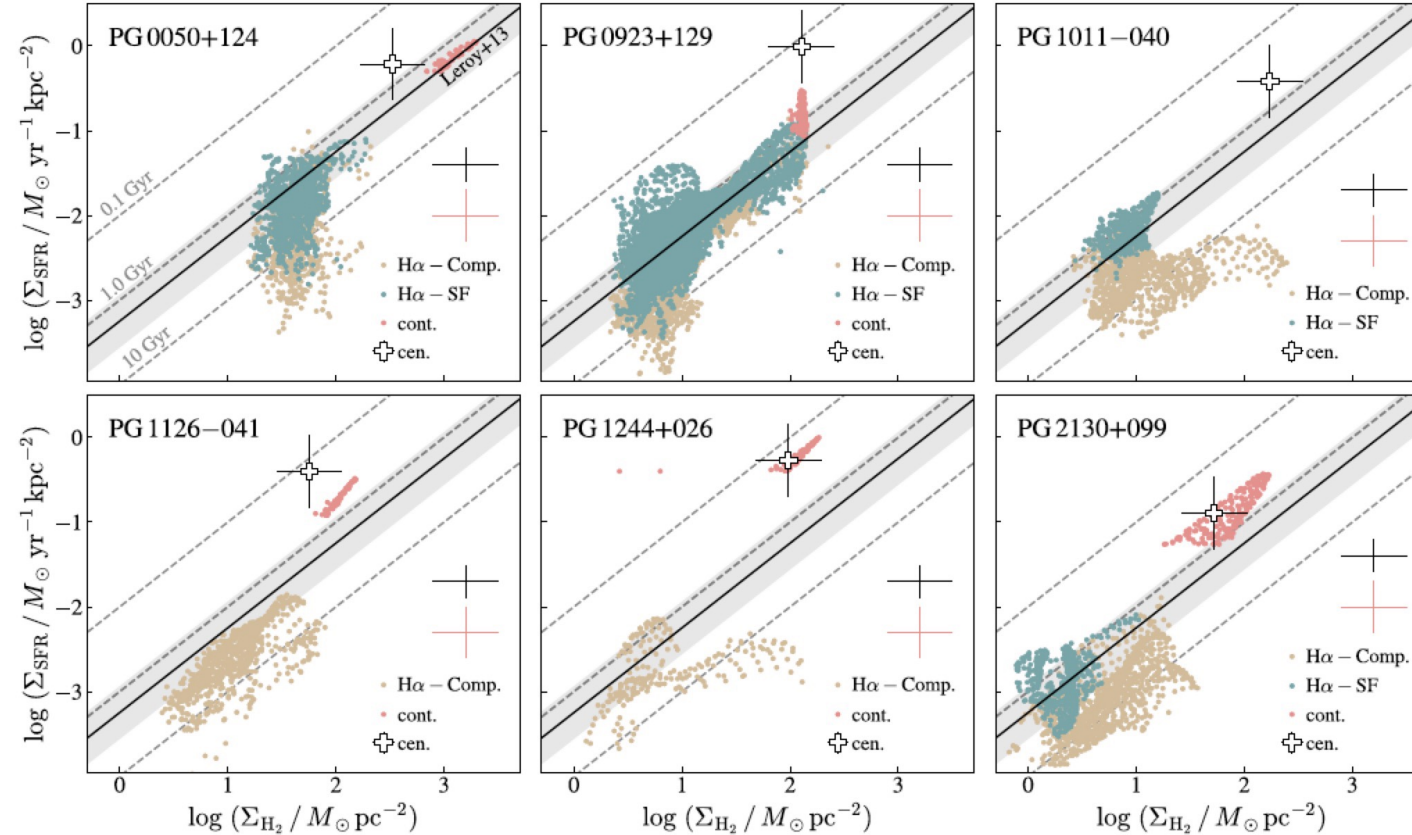
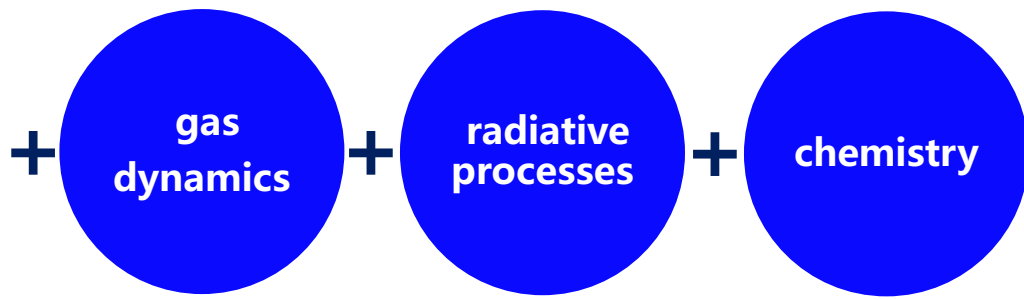


Figure 6. The PG quasar host galaxies in the Kennicutt–Schmidt diagram (Kennicutt 1998a). The red points show the Σ_{SFR} values estimated from the rest-frame 230 GHz continuum data. We show with green circles the MUSE pixels classified as star-forming, while the other pixels are represented by brown circles. The plus sign presents the average host galaxy central value computed over the host galaxy area where we miss the optical light emission owing to AGN deblending flux oversubtraction. The solid line shows the best-fit values of t_{dep} derived by Leroy et al. (2013) for local inactive galaxies, with the shaded region corresponding to the 1σ uncertainty. The dashed lines represent values of constant t_{dep} , as labeled in the top left panel. The regions with obscured and unobscured star formation activity are characterized by t_{dep} values similar to or slightly higher than those measured from the inactive systems, indicating starburst-like activity. The black and red colored error bars located above the legend correspond to the typical 1σ uncertainty of the H α - and continuum-based pixel-wise data.

Molina, Ho, et al 2023





gas process linked to stellar radiation

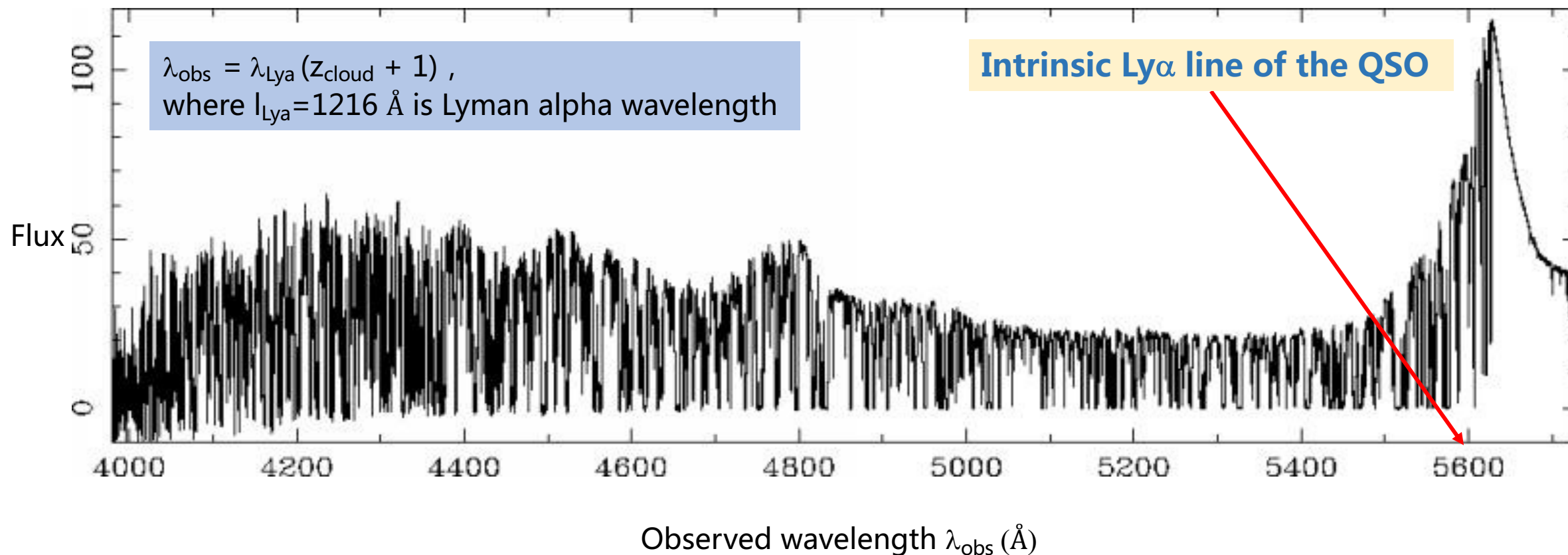
**One robust result from cosmological hydrodynamic simulations marks
outer gas boundary conditions due to photo-heating by stars & QSOs**

Lyman alpha forest



Numerous absorption lines were seen in the spectra of high- z quasars, hence the (Lyman alpha) forest

QSO at $z=3.6$



**Cosmological
significance**

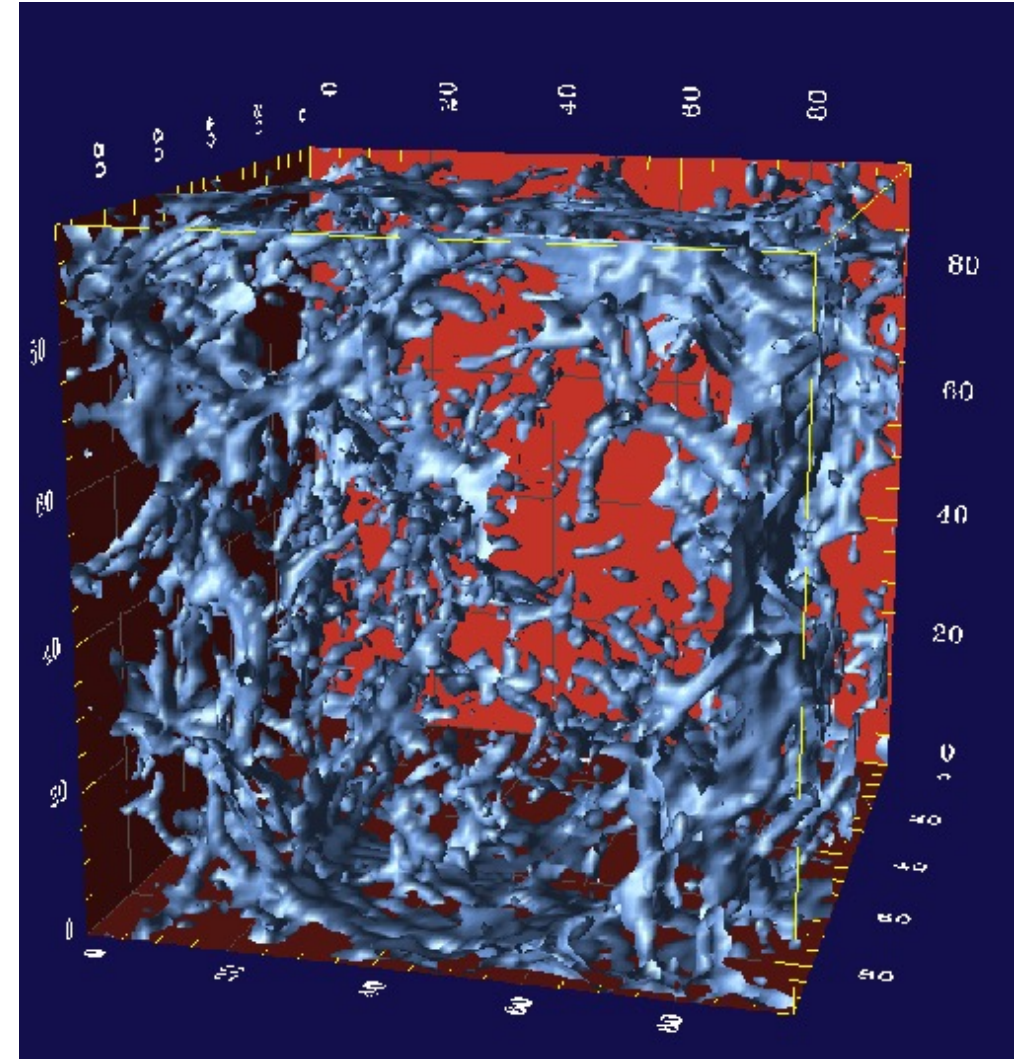
@ $z \geq 3$ the Lyman alpha forest contains
>95% of the baryons in the universe

**Heuristic
theoretical
models**

- Galactic outflows (Fransson & Epstein 1982)
- Confined by intergalactic hot gas (Sargent 1980; Ostriker & Ikeuchi 1983)
- Minihalos (Rees 1986; Ikeuchi 1986)
- Infalling gas to dark matter halos (Bond, Szalay & Silk 1988)
- Outer galactic disks (Charlton, Salpeter & Hogan 1993)
- Debris from merging satellites (Wang 1993; Morris & van den Bergh)
- Primordial linear density fluctuations (Bi 1993)
- Confined by hot gas in halos (Mo 1994)

Significance

- Lyman alpha forest absorption is the **fluctuating Gunn-Peterson optical depth field**, stemming from photo-ionized, nonlinear, large-scale baryonic structure of the universe at moderate redshift (Cen et al. 1994, ApJ, 437, L9)
- This theory quantitatively and systematically explains a whole array of the Lyman alpha forest observables (dn/dz , dn/dN_{HI} , b , T , ...)
- It is a **parameter-free theory** based solely on the Λ CDM cosmological model itself and little affected by uncertainties due to feedback processes



Isodensity surfaces @ $3 \times$ mean density @ $z=3$

Cen & Simcoe 1997, ApJ, 483, 8

One major contemporary method for measuring cosmological parameters

- **the parameter-free** Lyman alpha forest theory depends only on cosmological model parameters
- laid foundation for measurements of Lyman alpha forest flux distribution as one of **pillar methods** for determination of **cosmological parameters, neutrino mass, ...**

Examples: cosmological parameters and dark energy equation of state constraints

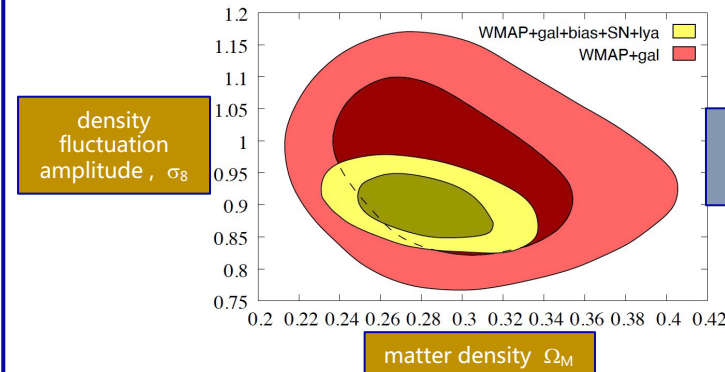


FIG. 5: 68% and 95% contours in the (Ω_m, σ_8) plane showing previous constraints from WMAP and galaxy clustering with the new data.

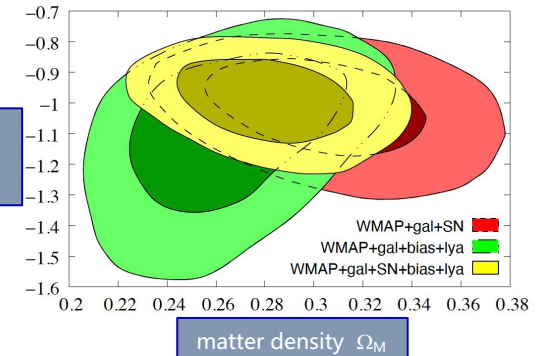
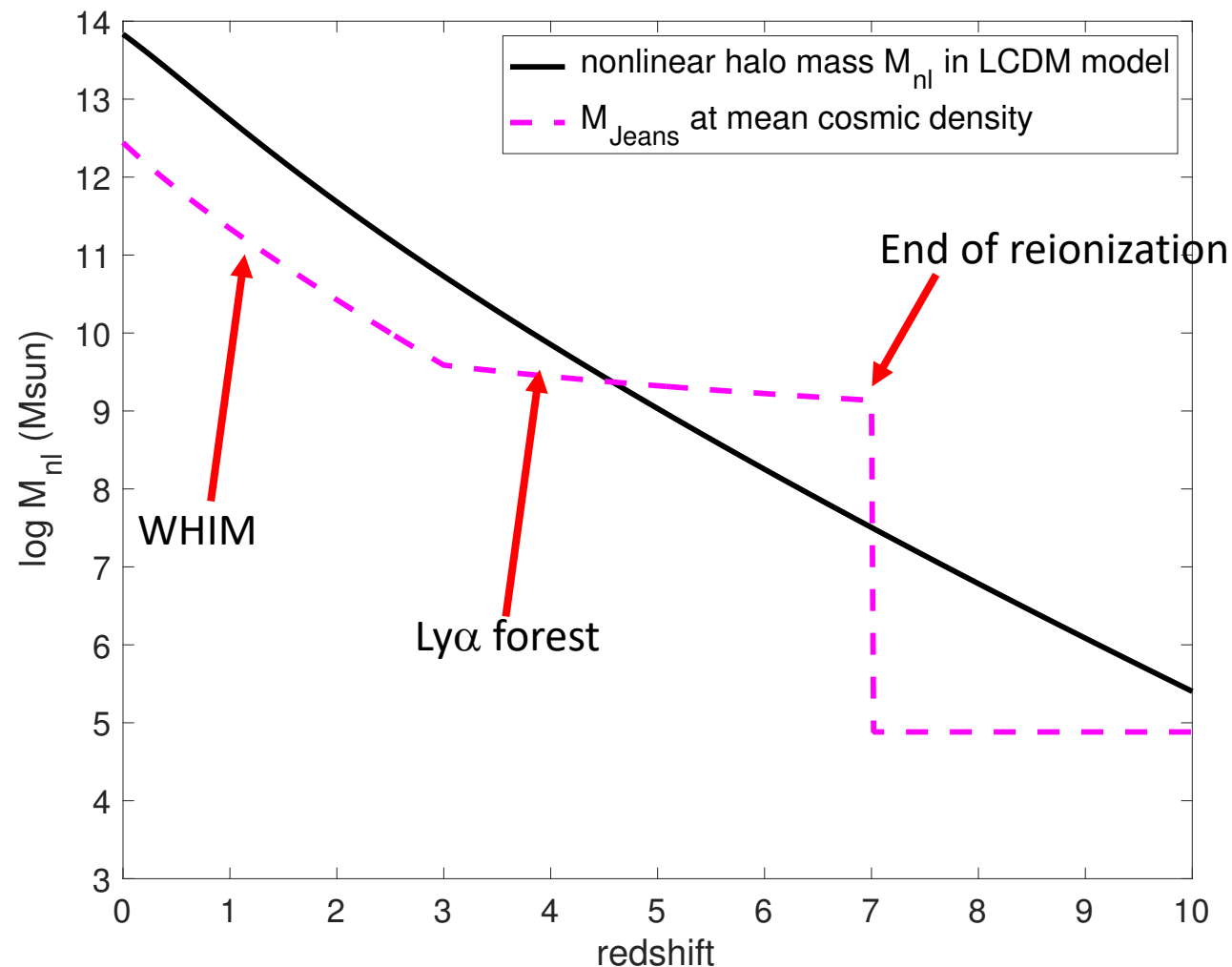


FIG. 7: 68% and 95% contours in the (Ω_m, w) plane showing previous constraints without SDSS-lya and bias, constraints without SNIa, and combined constraints. In all cases the data are consistent with a cosmological constant model ($w = -1$).

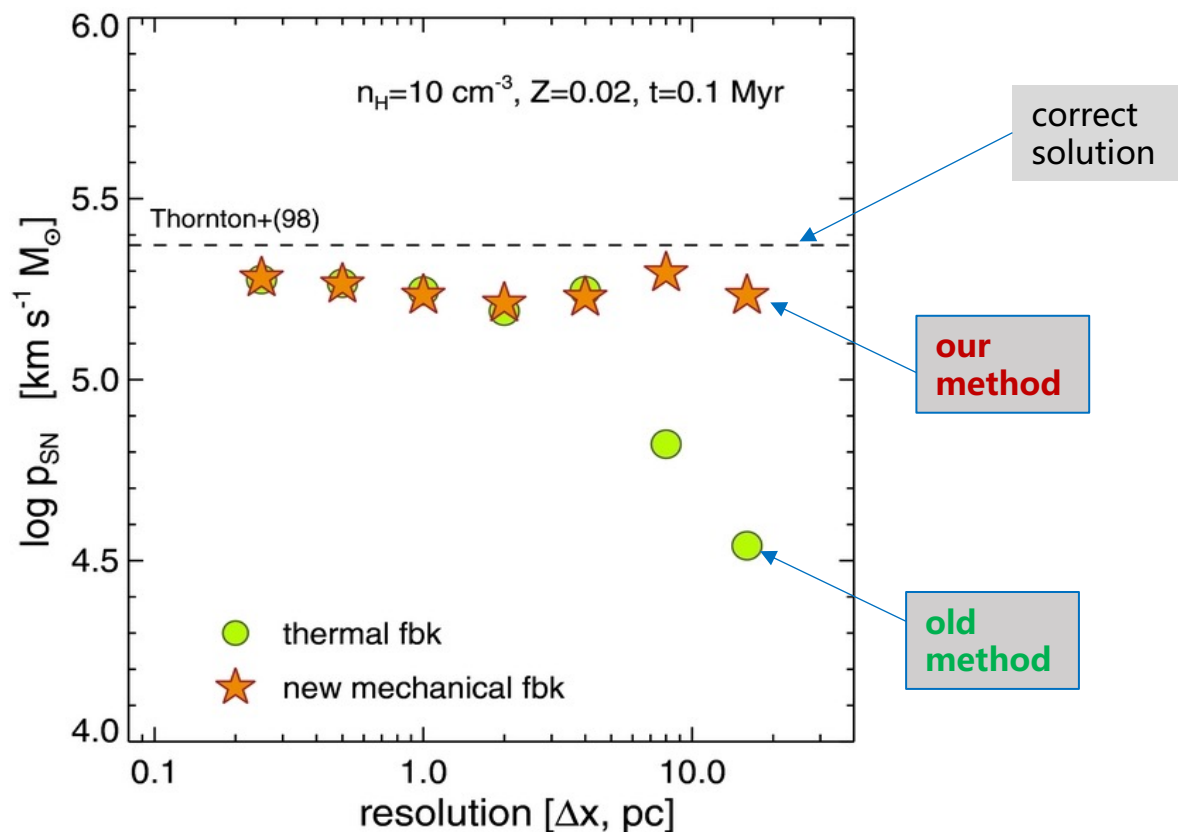
Seljak et al. 2005, PRD, 71, 103515;
see also Dawson et al. 2012, AJ, 145, 10
and Palanque-DeLabrouille et al. 2020, JCAP, 04, 038

cosmological origin of Ly α forest & discovery of warm-hot IGM

- Completion of reionization of the universe at $z \sim 7$
- photo-heated intergalactic medium imposes the Jeans mass scale for galaxy formation at $z=3-7$
- gravitational shock heated intergalactic medium sets the Jeans mass scale for galaxy formation at $z < 3$



a well known physical process, but an **effective** implementation with **minimum # of free parameters** was illusive until recently



Kimm & Cen 2014

From 1988, supernova feedback was added to simulations (Katz, Hernquist, Cen, ...) but ineffective

Kimm & Cen (2014) put forth a method to model supernova feedback, realizing that the Sedov-Taylor blastwave with cooling has in practice two distinct phases:

- **energy conserving phase:** $v = (2 E/m_{\text{shell}})^{1/2}$
- **momentum conserving phase:** $v = p_{\text{term}}/m_{\text{shell}}$
- **E conserving → p conserving has a rapid transition with a brief duration**
- **transition time is expressible by shell mass, which is a very weak function of (density, metallicity),**
 $m_{\text{shell}} \sim 100 m_{\text{ejecta}}$

Internal vs external feedback : Why external feedback is easier?

External feedback energetically more economical

- to blow away a baryon from inside a halo, you need about $\epsilon_{\text{int}} = m_p \sigma^2$ energy (σ is the velocity dispersion of the halo)
- to prevent a baryon from ever entering the same halo, you need $\epsilon_{\text{ext}} = \delta^{-2/3} \epsilon_{\text{int}} \approx 3\% \epsilon_{\text{int}}$ energy (δ is halo over-density) by heating it up at the mean cosmic density



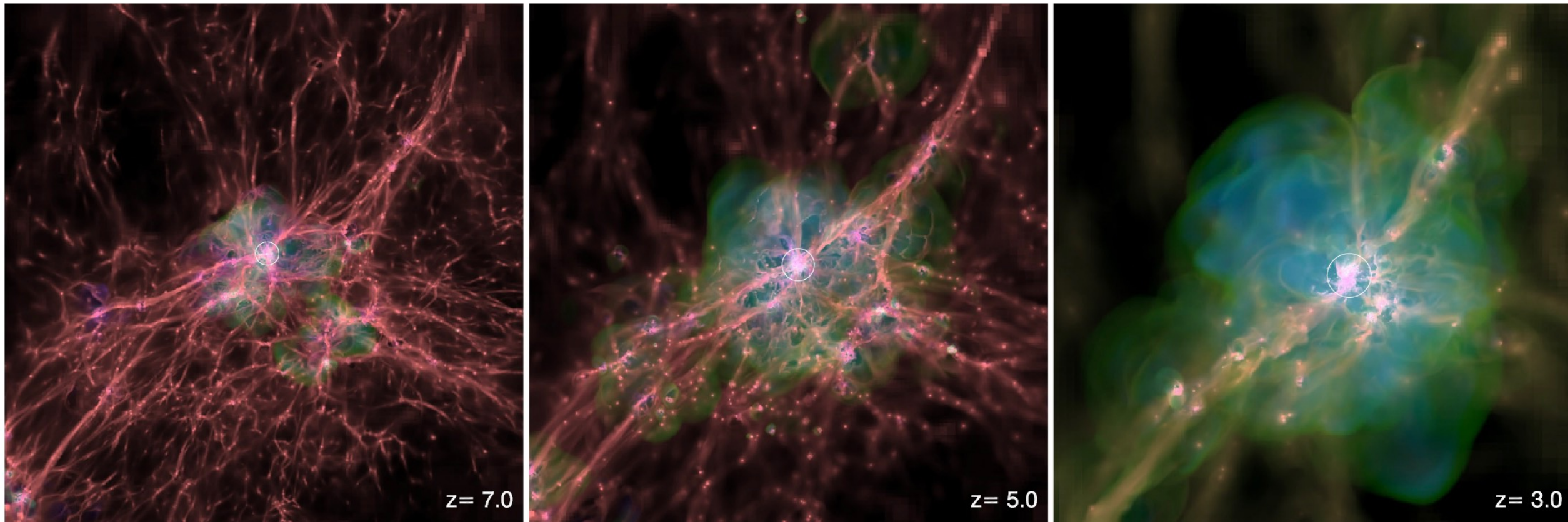
External SN feedback : **how important?**

It depends on how well, how much and how far supernova feedback energy can escape from galaxy

- SNe II feedback energy is tightly coupled to ISM; external feedback depends critically on the depth of gravitational potential well of the halo
- From dwarf galaxies, external SN feedback is likely very important, probably responsible for enriching IGM
- SNe Ia feedback energy is likely easier to escape due to weaker ISM coupling



External SN feedback : demonstration



Evolution of a galaxy of halo mass $10^{11.7}M_{\odot}$ at $z=0$, with correct implementation of supernova feedback

Red: density

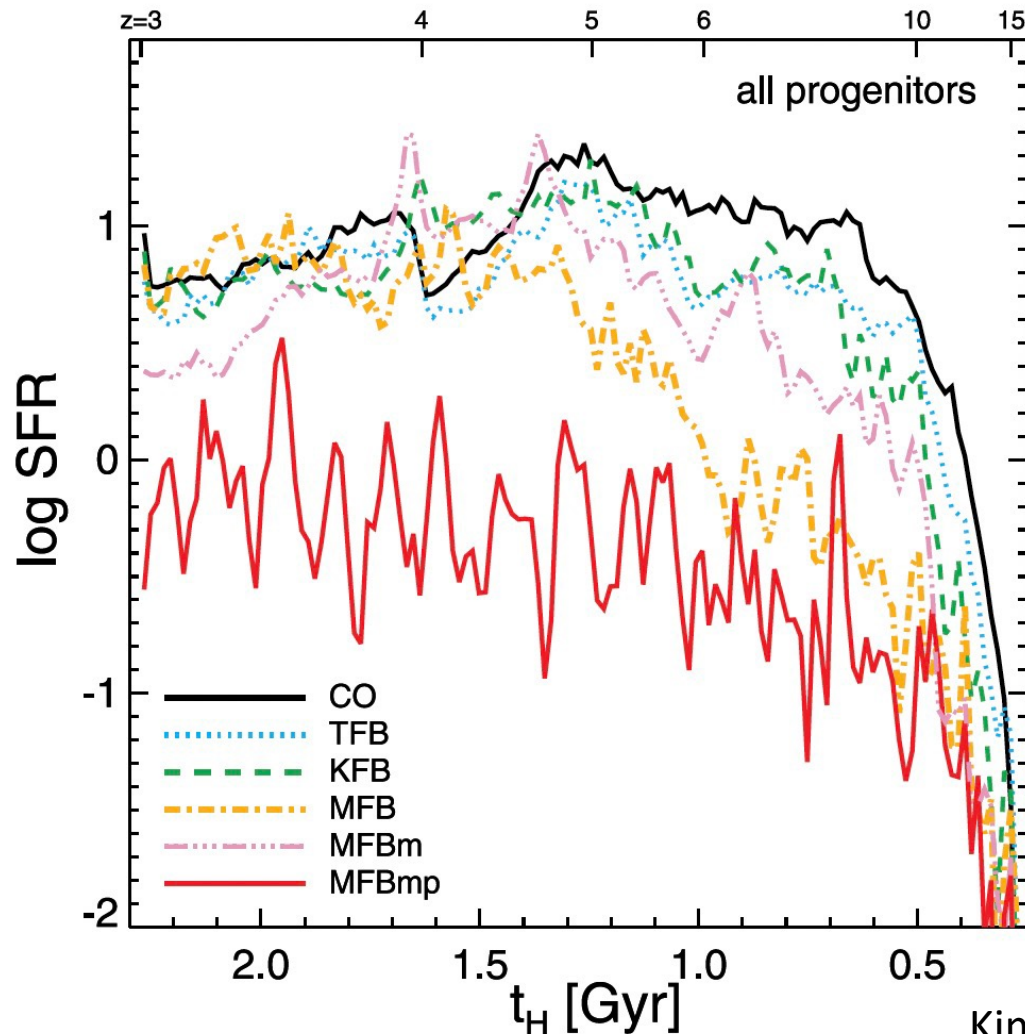
Green: metallicity

Kimm, Cen et al 2015

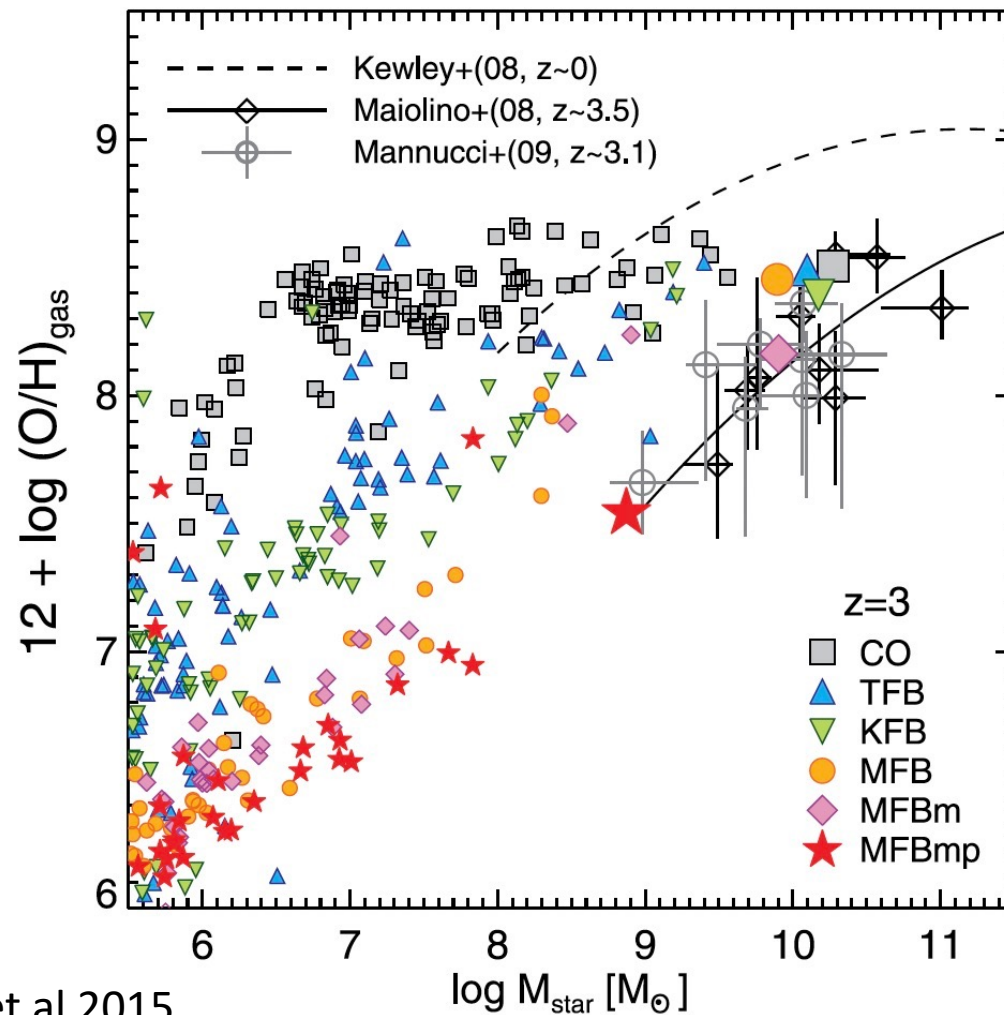


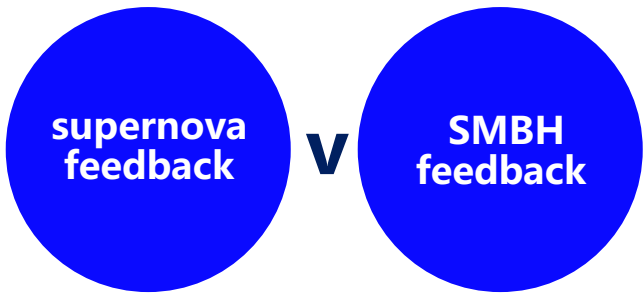
supernova
feedback

Galaxy formation outcomes depend **critically** on supernova feedback implementations



Kimm, Cen et al 2015





comparisons of energetics

THE ASTROPHYSICAL JOURNAL, 755:28 (14pp), 2012 August 10

CEN

Table 1
Comparison of SF and SMBH Energetics

No.	Form	SF	SMBH
(1)	Total radiation	$\epsilon_*(\text{rad}) = 7 \times 10^{-3}$	$\epsilon_{\text{BH}}(\text{rad}) = 2 \times 10^{-4}$
(2)	Ionizing radiation (≥ 13.6 eV)	$\epsilon_*(\text{LL}) = 1.4 \times 10^{-4}$	$\epsilon_{\text{BH}}(\text{LL}) = 3 \times 10^{-5}$
(3)	X-ray (2–10 keV)	$\epsilon_*(2\text{--}10\text{ keV}) = 9 \times 10^{-8}$	$\epsilon_{\text{BH}}(2\text{--}10\text{ keV}) = 5 \times 10^{-6}$
(4)	Mechanical	$\epsilon_*(\text{SN}) = 1 \times 10^{-5}$	$\epsilon_{\text{BH}}(\text{BAL}) = (0.2\text{--}2.8) \times 10^{-5}$
(5)	Radio jets	$\epsilon_*(\text{jet}) = 0$	$\epsilon_{\text{BH}}(\text{jet}) = 4 \times 10^{-5}$

Notes. Under the assumption that $M_{\text{BH}}:M_{\text{BG}} = 2:1000$, a Salpeter IMF for stars and a radiative efficiency of SMBH accretion of 10% (Yu & Tremaine 2002), and energy output from both SF and SMBH in various forms are listed: (1) total radiation energy, (2) ionizing radiation, (3) X-ray radiation in the 2–10 keV band, (4) mechanical energy, and (5) radio jets.

Cen 2012



External feedback : the critical entropy

gas cooling time exceeds Hubble time, if $S > S_{\text{crit}}$

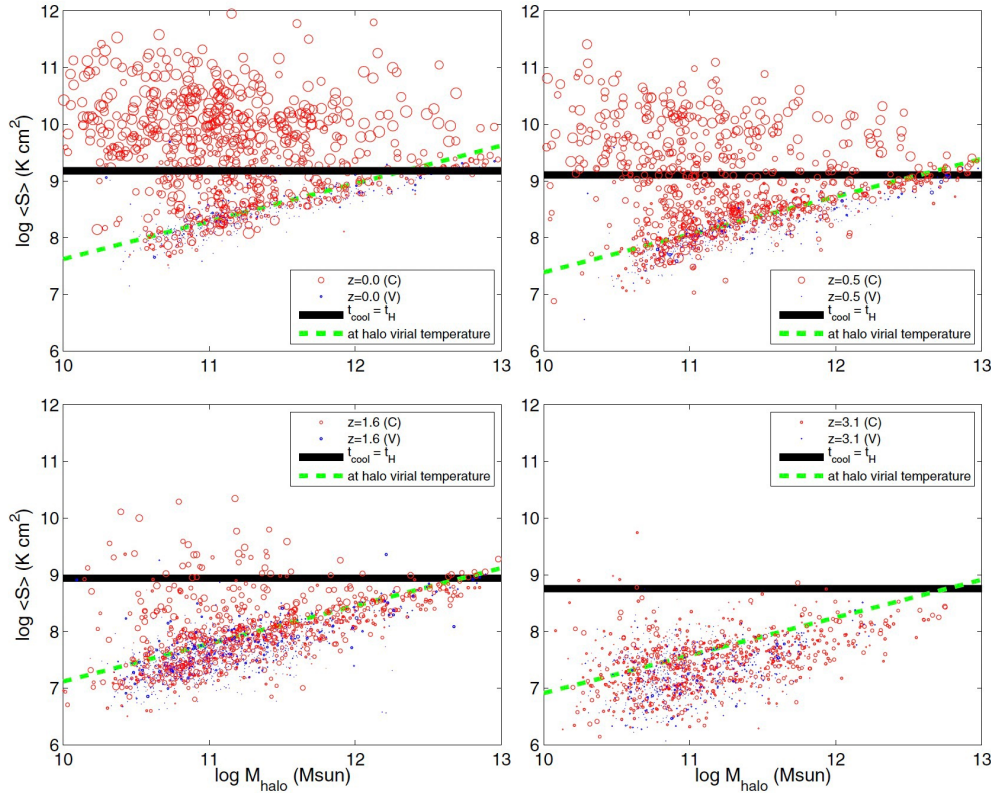
$$t_{\text{cool}} = \frac{(3/2)nk_B T}{n_e^2 \Lambda(T)} = S^{3/2} \left[\frac{3}{2} \left(\frac{\mu_e}{\mu} \right)^2 \frac{k_B}{T^{1/2} \Lambda(T)} \right]$$

$$\log[S_{\text{crit}}/(\text{K cm}^2)] = 9.183 - 0.167z + 0.0092z^2$$

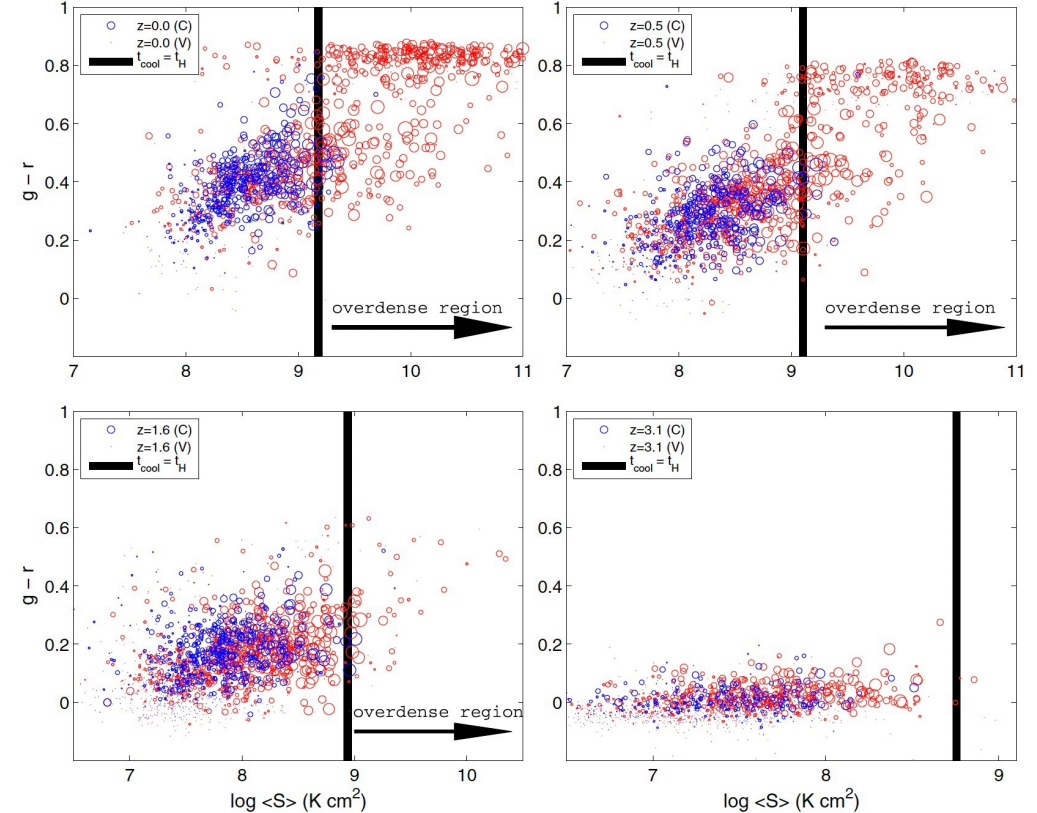
External feedback : the critical entropy

When environmental $S > S_{\text{crit}}$, galaxies turn red

THE ASTROPHYSICAL JOURNAL, 741:99 (15pp), 2011 November 10



THE ASTROPHYSICAL JOURNAL, 741:99 (15pp), 2011 November 10



CEN

[Submitted on 4 Apr 2023]

The CAMELS project: Expanding the galaxy formation model space with new ASTRID and 28-parameter TNG and SIMBA suites

Yueying Ni, Shy Genel, Daniel Anglés-Alcázar, Francisco Villaescusa-Navarro, Yongseok Jo, Simeon Bird, Tiziana Di Matteo, Rupert Croft, Nianyi Chen, Natalí S. M. de Santi, Matthew Gebhardt, Helen Shao, Shivam Pandey, Lars Hernquist, Romeel Dave

We present CAMELS-ASTRID, the third suite of hydrodynamical simulations in the Cosmology and Astrophysics with Machine Learning (CAMELS)

Current State of Affairs

- **Most simulations are focused on internal feedback from AGN**
- **AGN feedback have presently modelled with two modes**
 - (1) QSO mode (e.g., high energy flux isotropic winds)
 - (2) jets mode (e.g., low energy flux, long duration)
- **What is missing is external feedback from AGN, e.g., high energy FRI/FRII radio jets**
- **Inclusion of this key external feedback physics may alleviate the need of too many free parameters**

need to re-establish simulations as first-principle based

three significant challenges

- We need to cull an unbiased set of observational data on AGN that is impervious to personal amplifications, in
 - (1) winds in all forms
 - (2) radio jets (small and large)
- We then need to understand and extract tangible and implementable physics from such a data set
- We then need to combine focused simulations/theoretical works on AGN scales and make connecting simulations (e.g., between radio jets scale and galactic to extragalactic) to devise physics-based implementation schemes
- This is a difficult and major challenge, i.e., a great **opportunity for students** who are interested

Conclusions

- Through collective efforts from 1970s to 2010s, dark matter gravitational dynamics is essentially well understood, i.e., **we understand properties of dark matter halos in space and time** via simulations
- Through collective efforts from 1990 to now, we have made significant breakthroughs in understanding the thermodynamic properties of the **intergalactic gas in the cosmic web** at low and high redshift
- The core problem of galaxy formation has one last major hurdle: feedback processes from the growth of **supermassive black holes**
- Efforts to meet this challenge will be **richly rewarded physically** and will have the opportunity to have a say in the construction of the final theory of galaxy formation

